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SERMeQ model produces realistic retreat of 155 Greenland outlet glaciers

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Key Points:

- We simulated terminus advance and retreat of 155 ocean-terminating outlet glaciers that drain the Greenland Ice Sheet.
- Our simulated terminus positions lie within the observational range for 40% of observed terminus positions.
- Our model consistently overestimates retreat rates and ice mass loss, providing an upper bound for future sea level projections.

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Abstract

The rate of land ice loss due to iceberg calving is a key source of variability among model projections of 21st century sea level rise. In Greenland, where ice drains to the ocean through hundreds of outlet glaciers, it has been especially challenging to account for iceberg calving from glaciers smaller than typical model grid scale. Here, we apply an efficient, physically-based network flowline model (SERMeQ) forced by surface mass balance to simulate decadal terminus position change of 155 grounded outlet glaciers of the Greenland Ice Sheet—resolving five times as many outlets as was previously possible. We compare these simulations with observed changes in terminus position and find that SERMeQ produces generally realistic rates of retreat. Moreover, SERMeQ is designed to overestimate retreat and can be used to provide an upper bound on forward projections of the dynamic mass loss from the Greenland Ice Sheet associated with different climate projections.

1 Introduction

The Greenland Ice Sheet is currently the largest single contributor to global mean sea level rise (van den Broeke et al., 2017). It discharges ice mass to the ocean through three main processes: release of surface meltwater, submarine melting where ice is in contact with the ocean, and the detachment (calving) of icebergs. The ice mass lost to submarine melting has only recently been directly observed (Sutherland et al., 2019) and remains difficult to estimate for the whole ice sheet (Beckmann et al., 2018), but it is clear that enhanced surface melting and calving processes have resulted in increased mass discharge since the late 1990s (van den Broeke et al., 2016; Enderlin et al., 2014; Khan et al., 2014).

Processes that control surface melt are increasingly resolved in regional models (Mottram et al., 2017; Noël et al., 2018). Iceberg calving, by contrast, remains poorly understood, with multiple contradictory parameterizations incorporated into ice sheet/glacier models (Benn, Cowton, et al., 2017). Furthermore, iceberg calving can remove mass more rapidly than is possible through melting alone, contributing to rapid tidewater glacier retreat through mechanisms like tidewater glacier instability (Meier & Post, 1987) and the recently-described Marine Ice Cliff Instability (Bassis & Walker, 2012; Pollard et al., 2015).

Simulating discharge from the Greenland Ice Sheet is further complicated by the local factors affecting ice discharge at the nearly 200 outlet glaciers that connect the ice sheet to the ocean (e.g. Catania et al., 2018; Enderlin et al., 2018). For all but the largest outlets, iceberg calving occurs at smaller scales than are captured in continental-scale ice sheet models. Existing estimates of dynamic mass loss from Greenland outlets have come from extrapolating perturbations on the largest outlets (Price et al., 2011; Nick et al., 2013), simulating the sea level contribution from only selected outlets (Choi et al., 2017; Morlighem et al., 2019), or simulating the entire ice sheet at a spatial resolution of 500 m (Aschwanden et al., 2016, 2019). Despite these achievements, more than 100 outlet glaciers, responsible for $\sim 1/3$ of current Greenland Ice Sheet discharge (Enderlin et al., 2014), are not routinely simulated, and their dynamics cannot necessarily be inferred from the dynamics of larger outlets. Another layer of spatial complexity arises in that many outlet glaciers collect ice from several interacting tributary branches that are themselves also smaller than typical ice sheet model grid scale. The small scale of tributary glacier networks feeding outlets makes them especially challenging to simulate in continental ice sheet models, requiring model resolution of hundreds to tens of meters to adequately resolve.

A more fundamental challenge in projecting mass loss due to calving is the incompatibility of fracture-driven iceberg calving with the assumption of continuum deformation inherent in most ice sheet models (e.g. Price et al., 2015; Winkelmann et al., 2011;

Greve, 2000). Simple empirical parameterizations can relate calving rate to continuous variables, such as proglacial water depth (Brown et al., 1982; Hanson & Hooke, 2000), but may not hold into the future as climate forcing enters a new statistical regime. Physically-based calving laws, such as the fracture field approach developed by Albrecht and Levermann (2012) or von Mises calving law developed for Greenland by Morlighem et al. (2016), often impose an empirically-adjustable calving rate parameter. Recent work has sought to simulate ice failure using continuum damage mechanics, with some success in a variety of case studies (Borstad et al., 2012; Duddu et al., 2013; Krug et al., 2014; Sun et al., 2017; Mercenier et al., 2019). However, at present the evolution of the damage field through a damage production function is also empirical, with multiple tuned parameters that are poorly constrained by laboratory or field measurements (Emetç et al., 2018). Another recent approach couples a granular model that allows true fracture and calving to a finite-element model that solves an approximation to the Stokes equations for viscous deformation, offering a very promising basis for process-scale simulation of fully-dynamic calving (Benn, Åström, et al., 2017). Unfortunately, the coupled approach remains too computationally expensive for century-scale projections. Despite their promise, neither continuum damage models nor granular calving models have been able to reproduce observed multi-annual evolution of calving front positions in Greenland.

Improving projections of 21st-century sea level rise requires models that can (i) reproduce complex patterns of glacier advance and retreat currently observed in Greenland and (ii) efficiently simulate mass loss due to iceberg calving from individual outlet glaciers for a spectrum of climate scenarios. To address this, we have developed a simple model to simulate advance, retreat, and dynamic mass loss due to calving on networks of marine-terminating glaciers (Ultee & Bassis, 2016, 2017; Bassis & Ultee, 2019). Our model framework, called SERMeQ, manages computational expense so that it is possible to directly simulate decade-to-century-scale evolution of hundreds of outlet glaciers in response to surface mass balance forcing across multiple climate scenarios. This explicit simulation capability, together with recent observations of more than 200 Greenland outlet glacier termini (Joughin et al., 2015, updated 2017a), makes it possible to evaluate our model’s performance in a wide range of glacier environments. Here, we test its ability to reproduce present-day observed changes in terminus position of 155 Greenland outlet glaciers, providing one of the largest validations of any calving parameterization. We also demonstrate the calculation of equivalent sea-level contribution associated with the glacier retreat we simulate. On the basis of this validation, our model physics can be incorporated into global glacier and ice sheet models to compute a physically-consistent upper constraint on the century-scale glaciological contribution to global sea level rise.

2 Methods

2.1 SERMeQ ice dynamics model

SERMeQ—the Simple Estimator of Retreat Magnitude and ice flux (Q), *sermeq* meaning “glacier” in Greenlandic—is a width-averaged, vertically-integrated model that determines centerline glacier surface elevation corresponding to a given terminus position. The ice dynamics are based on a perfectly-plastic limiting case of a viscoplastic rheology (Nye, 1951; Bassis & Ultee, 2019), with modifications to allow calving at a grounded ice-water interface (Ultee & Bassis, 2016) and interaction between multiple tributary glaciers (Ultee & Bassis, 2017). Our flowline-modeling approach is compatible with other flowline-based models such as the Open Global Glacier Model (Maussion et al., 2019), but SERMeQ focuses specifically on near-terminus dynamics of marine glaciers.

Rather than imposing an empirical calving rate, SERMeQ self-consistently calculates the maximum rate of terminus advance or retreat at each time step for a given climate forcing. Terminus position evolves in response to near-terminus stretching, bedrock

116 topography, and changes in catchment-wide surface mass balance as described in Ultee
117 (2018) and Bassis and Ultee (2019),

$$\frac{dL}{dt} = \frac{\dot{a} - H \frac{\partial U}{\partial x} - U \frac{\partial H}{\partial x}}{\frac{\partial H_y}{\partial x} - \frac{\partial H}{\partial x}}. \quad (1)$$

118 In Equation 1, $H = H(x, t)$ is the ice thickness, $U = U(x, t)$ the ice velocity, $\dot{a} = \dot{a}(x, t)$
119 the net ice accumulation rate, H_y the thickness at which effective stress within the ice
120 reaches its yield strength (Equation S1), and all terms are evaluated at the instantaneous
121 terminus position, $x = L(t)$ (see Supplementary Text S1-2). For a change in terminus
122 position determined from Equation 1, SERMeQ calculates a new steady-state glacier sur-
123 face elevation profile (Supplementary Figure S1) and calculates change in glacier volume
124 above buoyancy. The latter produces a net contribution to global mean sea level (Fig-
125 ure 4, below).

126 The only adjustable model parameters are ice temperature T , which is used to calcu-
127 late the horizontal stretching rate $\partial U / \partial x$ at the terminus, and yield strength τ_y , which
128 is used to calculate the yield thickness H_y (Supplementary Text S1-S2). Both are ma-
129 terial quantities that can be independently constrained by laboratory and field measure-
130 ments. Crucially, we do not tune either of our parameters to match changes in termi-
131 nus position. Comparison of simulated with observed terminus position provides a com-
132 pletely independent validation.

133 Here, we extend the physical realism and applicability of our model to demonstrate
134 that it can be applied to simulate advance and retreat of a wide variety of calving glaciers.
135 Novel elements of SERMeQ specific to this application include upstream forcing with
136 surface mass balance from a regional climate model (Mottram et al., 2018), the auto-
137 matic selection of networks of flowlines with varying width (traced from Joughin et al.,
138 2015, updated 2017b), and the calculation of net sea level contribution associated with
139 changes in glacier terminus position.

140 2.2 Identification of flowline networks

141 We first identified 181 Greenland outlet glaciers that have multiple terminus po-
142 sitions recorded in Joughin et al. (2015, updated 2017a). For each glacier, we then de-
143 fined a network of interacting flowlines with spatially variable width by tracing ice sur-
144 face velocity from Joughin et al. (2015, updated 2017b). We extracted ice surface and
145 bed elevation from BedMachine version 3 (Morlighem et al., 2017) and applied a Gaus-
146 sian filter to produce width-averaged topography. Where the data suggested the pres-
147 ence of short, transient ice tongues, we removed the floating portion from consideration
148 and simulated the grounding line as the “terminus”. We removed three glaciers with long,
149 persistent ice tongues, as SERMeQ is unable to simulate their dynamics. Thirteen of the
150 181 outlets had initial termini grounded above sea level and iceberg calving is thus un-
151 likely to dominate dynamic mass changes there. We removed those thirteen glaciers from
152 consideration as well. Noisy or missing data that produced unphysical bed topography
153 caused us to remove ten additional outlets, leaving 155 glaciers for our analysis.

154 For the remaining 155 outlet glaciers, we defined the initial terminus as the grounded-
155 ice point along our central flowline that lies closest to the centroid of the 2006 terminus
156 reported in Joughin et al. (2015, updated 2017a). We optimized a single parameter, the
157 yield strength of ice, to best fit the 2006 observed surface profile, as described in Ultee
158 and Bassis (2017). We used a best-guess ice temperature T of -10° C and did not ad-
159 just it here. We then found the catchment-wide, annual mean surface mass balance forc-
160 ing for each outlet, \dot{a} in Equation 1, from HIRHAM regional climate model reanalysis
161 (Mottram et al., 2018; Rae et al., 2012; Lucas-Picher et al., 2012), and simulated result-
162 ing changes between 2006 and 2014 in ice extent (Figures 1-3) and volume above buoy-
163 ancy (Figure 4 and Supplementary Figure 1). Finally, we compared the simulated changes

164 in terminus position with observed changes reported in Joughin et al. (2015, updated
 165 2017a) for the same period. Because our optimization of τ_y considers only the initial ob-
 166 served surface profile, and the changes in terminus position are an independent response
 167 to changes in surface mass balance, this comparison examines an independent model pre-
 168 diction that is not tuned to match observations.

169 Our automated flowline extraction defined the model domain to end at the initial
 170 terminus position. On 20 of the 155 glaciers we studied, Equation 1 predicted a positive
 171 rate of length change that would take the terminus outside of the defined domain dur-
 172 ing the simulation. For those glaciers, we artificially projected flowlines seaward to al-
 173 low glacier advance.

174 **2.3 Comparison with observations**

175 We extracted all available terminus position records from (Joughin et al., 2015, up-
 176 dated 2017a) for each year within our simulated period: 2006, 2007, 2009, 2013, and 2014.
 177 Each terminus position record consists of one or more points; records with multiple points
 178 trace across-flow variation in terminus position. We projected all available points from
 179 a given record onto the central flowline of the corresponding glacier network, and we iden-
 180 tified the space between the most seaward and most landward points of that projection
 181 as the “observational range”. We also tracked the change over time in the position of
 182 the terminus centroid projected on the flowline, which we identified as the “observed (terminus-
 183 centroid) retreat rate”. Finally, we compared the simulated retreat rates with the ob-
 184 served terminus-centroid retreat rates (Figure 2a) and the simulated terminus positions
 185 with the observational range (Figures 2b-3).

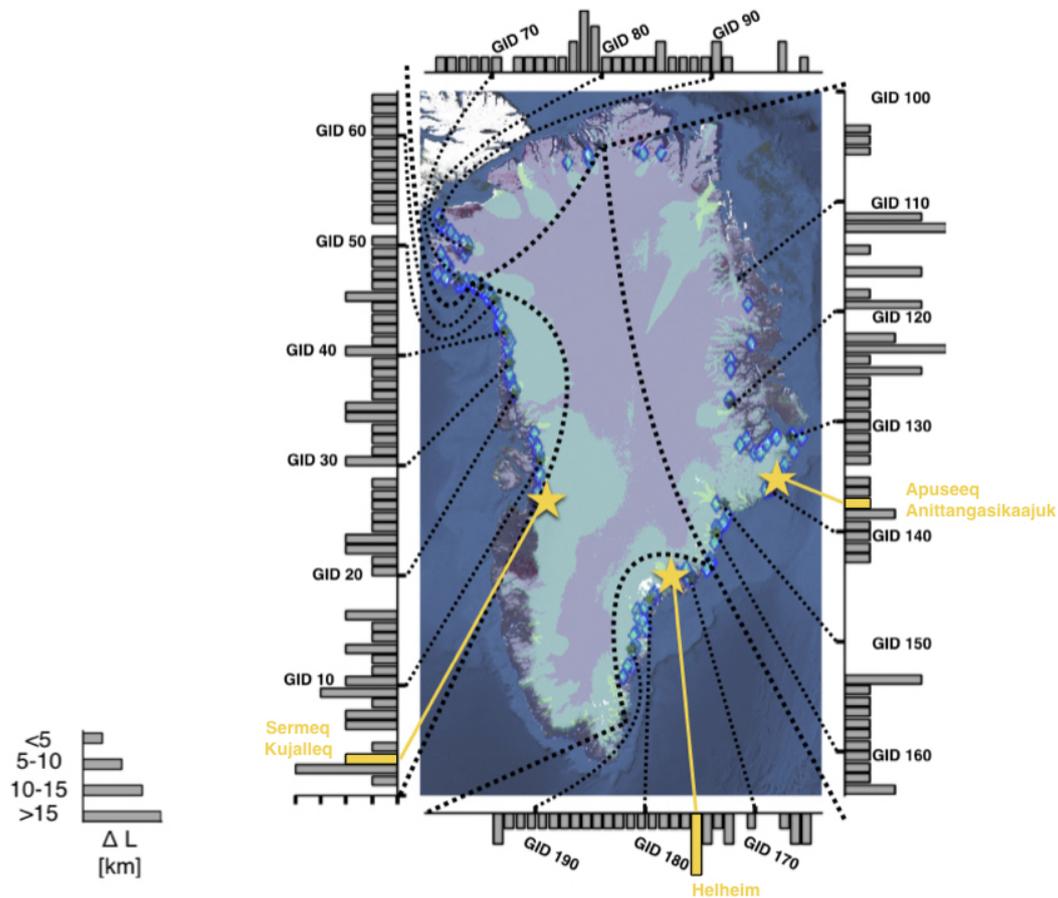
186 **3 Results**

189 **3.1 Realistic upper-bound retreat with only two physical parameters**

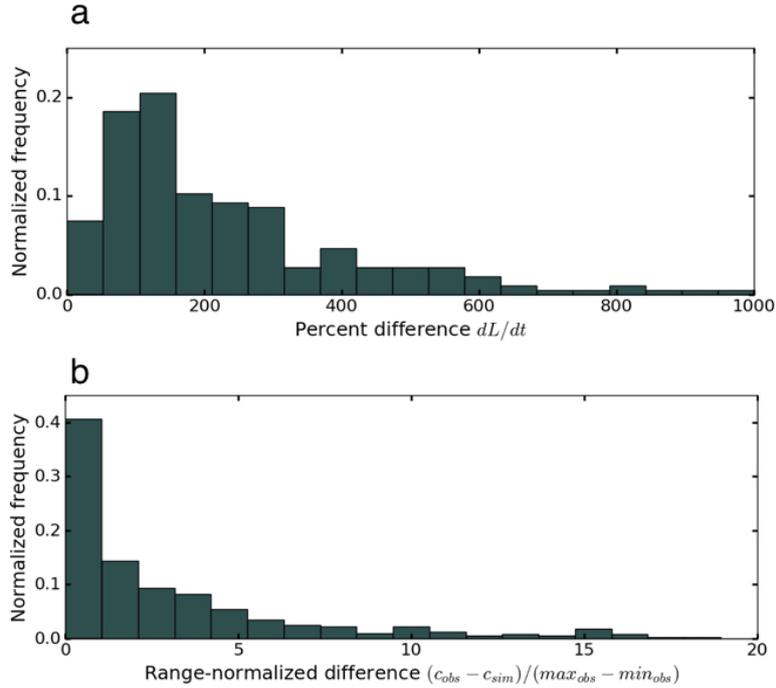
190 Figure 1 shows the total retreat we simulated for each glacier between 2006 and
 191 2014, arranged by approximate outlet position. SERMeQ simulates less than 5 km of
 192 length change during the observed period on most outlets. There is no relationship be-
 193 tween outlet glacier latitude and magnitude of simulated retreat: simulated glacier re-
 194 sponse to downscaled climate reanalysis forcing is not a simple function of annual av-
 195 erage temperature. Dynamic glacier response depends on glacier geometry, as previous
 196 studies have also highlighted (Felixson et al., 2017; Catania et al., 2018).

197 Equation 1 includes an assumption that the glacier calving front is a yield surface,
 198 which produces a theoretical upper bound on calving retreat for a given glacier geom-
 199 etry and surface mass balance (see Bassis & Ultee, 2019). Thus, provided there are no
 200 significant errors in the bed geometry and surface mass balance used, we anticipate that
 201 SERMeQ-simulated rates of retreat will generally overestimate observed rates. Figure
 202 2 shows that SERMeQ satisfies this expectation and overestimates retreat. Figure 2a shows
 203 that 55% of simulated retreat rates are within a factor of two of the corresponding ob-
 204 served rate, with a long tail of overestimates up to a factor of 10 (1000% difference). Sim-
 205 ilarly, Figure 2b shows that 40% of simulated terminus positions are within the obser-
 206 vational range of the corresponding MEaSUREs terminus position, with a long tail of
 207 estimates falling outside the observational range. The long tails in Figure 2a-b illustrate
 208 that the theoretical maximum retreat rate of Equation 1 can far exceed the observed rate,
 209 for reasons we address in the Discussion and Supplementary Text S5.

210 The bulk model results shown in Figures 1 and 2 summarize multi-annual histo-
 211 ries of terminus position change. Figure 3a shows full histories of terminus position change
 212 for all glaciers simulated, including the 20 glaciers for which we artificially extended flow-
 213 lines to allow advance (indicated as “flowlines projected seaward” purple rectangles; see
 214 215 216 217 218 219 220 221 222 223 224 225 226 227 228



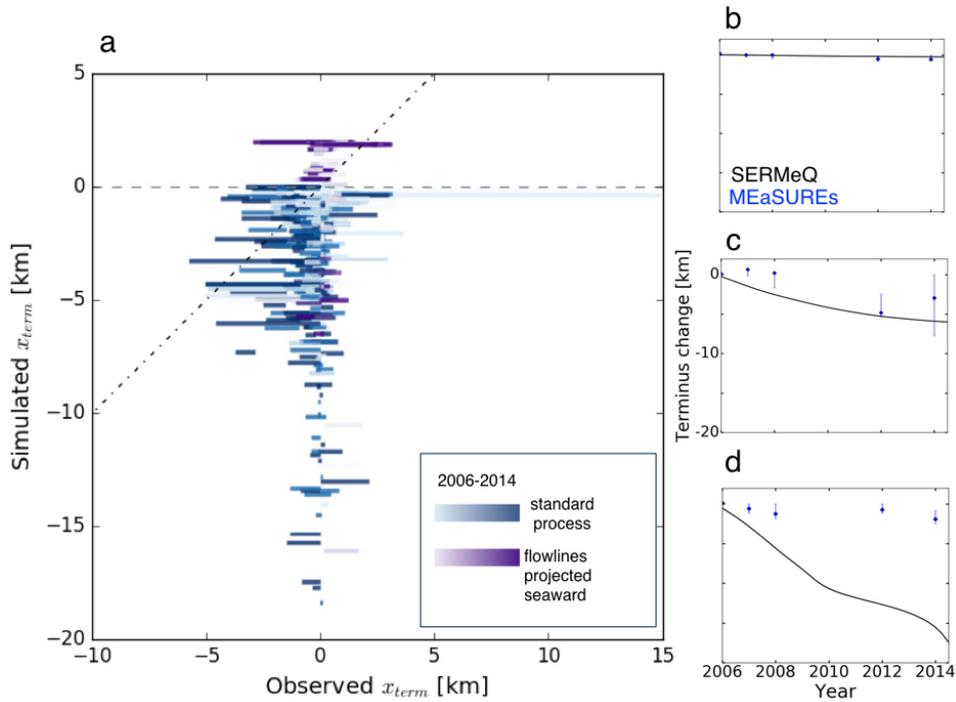
187 **Figure 1.** Map view of the 2006-2014 retreat simulated in this work. Bars indicate magnitude
 188 of simulated retreat for each glacier, with glaciers identified and ordered by their MEASURE
 189 outlet glacier ID number (1-200). Glacier ID 1, which is in the Disko Bay region, appears in
 190 the lower left; glacier IDs increase clockwise around the map border. Blue diamonds mark the
 191 map location of each outlet we simulated, and every 10th glacier ID is labelled and connected
 192 to its outlet location in black. A table of MEASURE glacier IDs and names appears in the
 193 Supplementary Material. Border spaces with no bar correspond to outlets where data was not
 194 sufficient to initialize a SERMeQ simulation, or where our analysis indicated SERMeQ would not
 195 be applicable (see Section 2). Yellow stars and bars show the case-study glaciers highlighted in
 196 Figure 3. Coloured overlay on the satellite map is ice velocity derived from Sentinel-1 observa-
 197 tions (ENVEO, 2017), shown on a logarithmic scale such that fast-moving outlet networks appear
 198 brighter than slow-moving inland ice.



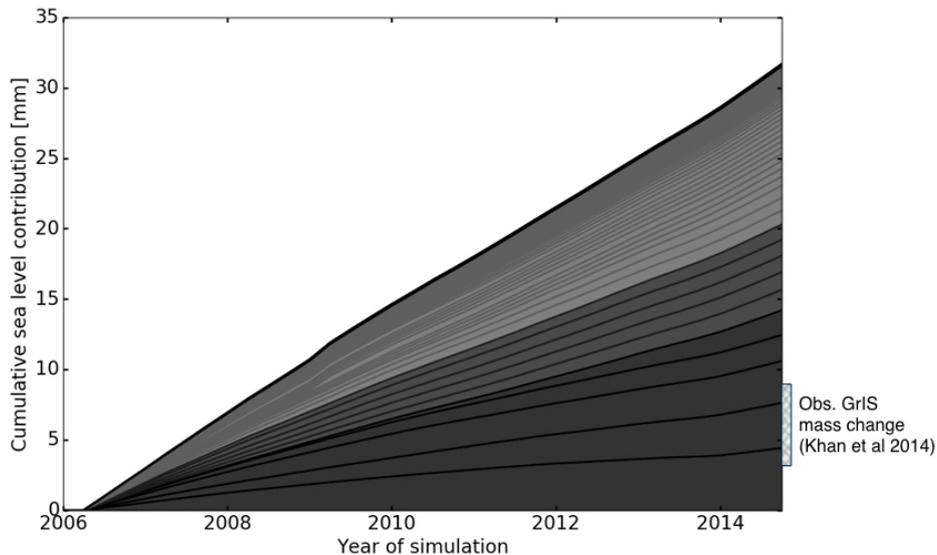
220 **Figure 2.** Histogram of differences in observed vs. simulated length change in Greenland out-
 221 let glaciers over the period 2006-2014. Panel (a) Percent error in rate of length change, dL/dt ,
 222 the quantity computed by Equation 1; Panel (b) Range-normalized difference in terminus posi-
 223 tion, where values ≤ 1 indicate a simulated terminus position within the observed range for a
 224 given year.

229 Section 2.2). Figure 3b-d compare observed and simulated terminus position change for
 230 example glaciers where SERMeQ underestimates, correctly captures, or overestimates
 231 the observed rate of retreat. Apuseeq Anittangasikkaajuk, which is 2 km wide at the ter-
 232 minus and has a small floating ice tongue, is one of a handful of outlets where SERMeQ
 233 underestimates observed retreat. The simulated terminus positions are still within the
 234 (small) observational range in that case. On Sermeq Kujalleq (Danish: Jakobshavn Is-
 235 bræ), a very large and well-studied outlet glacier on the southwest coast of Greenland,
 236 the simulated retreat of 6 km is comparable to observed retreat. SERMeQ overestimates
 237 retreat of Helheim Glacier, a large and high-flux glacier on Greenland’s east coast whose
 238 terminus approaches flotation. We address factors contributing to these estimates be-
 239 low and in Supplementary Text S5.

253 We achieve the realistic upper-bound simulations presented here with only two ad-
 254 justable model parameters: the yield strength of glacier ice τ_y and ice temperature T .
 255 Both are physical quantities constrained by laboratory and field observations. The best-
 256 fit yield strengths we find for the Greenland outlet glaciers simulated here range from
 257 55-200 kPa, well within the range of 50-500 kPa suggested by previous works (Nimmo,
 258 2004; O’Neel et al., 2005; Cuffey & Paterson, 2010). We use an ice temperature of -10°C ,
 259 which is within the range expected from simple physical scaling (van der Veen, 2013),
 260 observations (Clow et al., 1996), and modeling (Greuell & Konzelmann, 1994). Neither
 261 τ_y nor T is adjusted during simulations. Our simulated rate of terminus retreat/advance
 262 emerges as a dynamic glacier response to climate forcing and glacier geometry and does
 263 not rely on any tuning to match observations. It is possible an improved match to ob-
 264 served retreat rates could be found if we did allow parameters to vary within glacier catch-



240 **Figure 3.** Comparisons of observed and simulated terminus position change for (a) all glaciers
 241 simulated, with blue dashed 1:1 line; (b) Apuseeq Anittangasikkaajuk (glacier ID 137), where
 242 SERMeQ underestimates the true rate of retreat; (c) Sermeq Kujalleq (glacier ID 3), where
 243 SERMeQ captures observed retreat; (d) Helheim Glacier (glacier ID 175), where SERMeQ
 244 overestimates retreat. In panel (a), blue shaded rectangles show all observed-simulated ter-
 245 minus pairs for glaciers under fully automated processing, and purple shaded rectangles show
 246 all observed-simulated terminus pairs for glaciers that required artificial extension of flowlines.
 247 Rectangle width shows observational range and shade indicates time, with light colors in 2006
 248 progressing to dark colors in 2014. In panels (b-d), black curves indicate SERMeQ-simulated
 249 terminus positions, while blue markers indicate MEaSUREs observations. The blue lines show
 250 the most-advanced and most-retreated parts of the terminus projected onto the centerline, and
 251 blue diamonds indicate the centroid of the observed terminus projected onto the centerline. Plots
 252 share both x- and y-axis scales.



268 **Figure 4.** Cumulative ice dynamic contribution to global mean sea level due to SERMeQ-
 269 simulated upper-bound retreat in the period 2006-2014. Each curve indicates a contribution from
 270 an individual glacier network, and the uppermost curve is the cumulative ice-dynamic sea level
 271 contribution from all outlets. The grey hatched rectangle indicates total Greenland Ice Sheet
 272 mass loss over the period 2006-2012, extrapolated to 2014, based on observations by the GRACE
 273 satellite (Khan et al., 2014).

265 ments or over time. However, this would introduce empirical tuning that cannot be in-
 266 dependently constrained by laboratory or field observations.

267 3.2 Sea level contribution

274 As each glacier’s terminus position changes according to Equation 1, SERMeQ ad-
 275 justs the upstream surface profiles of all connected flowlines (Supplementary Figure S1).
 276 Assuming that centerline changes in ice surface elevation reflect width-averaged changes
 277 on each glacier branch, we can find the integrated change in ice volume above buoyancy
 278 and equivalent contribution to global mean sea level for each glacier over the observa-
 279 tional period. Figure 4 shows the SERMeQ-estimated cumulative sea level contribution
 280 over the 2006-2014 period from all of the 155 glaciers we simulated. The SERMeQ es-
 281 timate includes ice mass lost to calving as well as upstream dynamic drawdown, but does
 282 not include changes in ice mass due to changing surface mass balance.

283 We have constructed our model to produce an upper bound on glacier terminus re-
 284 treat rate, and the sea level contribution we compute is therefore an overestimate. Our
 285 goal is not to perfectly capture observed dynamic mass loss, but to illustrate the poten-
 286 tial for SERMeQ to compute and refine dynamic mass loss within large-scale simulations.
 287 Figure 4 shows that our estimate exceeds observational estimates for the same period
 288 by a factor of five. Given that we have made no tuning adjustments to the calving rate
 289 in this application, model agreement with observations is encouraging.

4 Discussion

The tendency to overestimate retreat supports the utility of our model for producing upper bounds on calving retreat and dynamic mass loss. In contrast to existing estimates of 21st-century calving loss, our approach does not impose a calving rate or outlet glacier speedup factor (DeConto & Pollard, 2016; Goelzer et al., 2013; Graverson et al., 2011; Pfeffer et al., 2008); instead, we calculate a theoretical maximum rate of calving retreat that can vary by glacier (Bassis & Ultee, 2019). Further, our model tracks terminus retreat and mass loss from multiple interacting branches of a glacier tributary network (Ultee & Bassis, 2017), ensuring that potentially important contributions are not overlooked. Within ice-sheet-scale models, our method could be implemented as a calving criterion at grounded ice-ocean interface cells or used as a module to enhance resolution of outlet glacier networks.

There are three notable sources of discrepancy between the modelled and observed retreat rates shown in Figures 2-3: (1) quality of available model input data, (2) performance of automated flowline selection algorithm, and (3) presence of floating ice. First, on small outlets that are rarely visited or studied in detail, the bed topography and climate reanalysis data used as input for SERMeQ may be poorly constrained. As a result, the simulated glacier evolves in response to conditions that do not accurately reflect the local environment, and the simulated change in terminus position is more likely to be inaccurate. Second, on small or slow-moving outlets, or where there are gaps in Sentinel-1 velocity data, our method for tracing flowlines is prone to error. As a result, the simulated glacier has unrealistic geometry and may flow over bedrock features that are not present in a true central flowline of the outlet. Finally, where floating tongues are present, we remove them and simulate the first grounded grid point as the “terminus”. This can change the near-terminus stress state, in some cases exposing an unstable wall of thick ice and initiating rapid retreat. Effects (1) and (2) are likely responsible for the underestimated retreat of Apuseeq Anittangasikkaajuk; effect (3) is likely responsible for the overestimated retreat of Helheim Glacier (see Supplementary Text S5). The first two effects can be mitigated with improved observational data and manual data processing where possible. The third effect reflects upper-bound retreat dynamics that are currently held in check by floating ice, but which we speculate could be activated if that floating ice were removed.

The current version of SERMeQ does not explicitly simulate frontal ablation by submarine melting, which can be a large component of mass loss from both floating tongues and grounded glacier fronts (Rignot et al., 2010; Enderlin & Howat, 2013; Wood et al., 2018). Our derivation of Equation 1, which we emphasise is an upper bound on retreat rate, is consistent with high submarine melt that prevents the glacier terminus from advancing (see Supplementary Text S4 and Ma, 2018; Ma & Bassis, 2019). However, changes in ocean conditions over time can affect glacier terminus dynamics such that the rate of terminus position change becomes closer to or farther from the theoretical maximum. For example, a decrease in submarine melt rate has been implicated in the recent slowing of Sermeq Kujalleq’s retreat (Khazendar et al., 2019). Future implementations of our method in larger-scale models may therefore benefit from modifications to account for time-varying submarine melt rates.

5 Conclusions

We have applied a flowline network model of ice dynamics, SERMeQ, to simulate an upper bound on annual to decadal-scale calving retreat of 155 Greenland outlet glaciers in response to variable climate forcing. Comparison with nearly a decade of terminus position records from MEaSUREs (Joughin et al., 2015, updated 2017a) shows that 55% of simulated retreat rates are within a factor of two of the observed rate. SERMeQ also evolves upstream surface elevation with each change in terminus position and computes

341 the resultant loss of ice mass above buoyancy. The model tends to overestimate retreat
 342 and will tend to overestimate the corresponding loss of grounded ice mass. The overes-
 343 timations of SERMeQ are consistent with efforts to find an upper bound on the ice-dynamics
 344 contribution to 21st century sea level rise. Our approach is especially promising in con-
 345 straining the dynamic sea level contribution from smaller outlet glaciers that are diffi-
 346 cult to resolve in larger-scale continental ice sheet models.

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 349 the MEASUREs project (Joughin et al., 2015, updated 2017a, 2010, 2015, updated 2017b),
 350 available from the National Snow and Ice Data Center. Surface mass balance forcing comes
 351 from the HIRHAM regional climate model for Greenland, maintained by the Danish Me-
 352 teorological Institute and available from [http://prudence.dmi.dk/data/temp/RUM/HIRHAM/
 353 GREENLAND/](http://prudence.dmi.dk/data/temp/RUM/HIRHAM/GREENLAND/). Python code for data processing (inc. network selection), simulation, and
 354 analysis is maintained in a public GitHub repository, which can be inspected at [http://
 355 github.com/ehultee/plastic-networks](http://github.com/ehultee/plastic-networks). This work is supported by the DOMINOES
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358 The authors have declared that no conflict of interest exists.

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