

Coversheet for “Intraseasonal Sea-Level Variability in the Persian Gulf”

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ABSTRACT

Satellite observations are used to establish the dominant magnitudes, scales, and mechanisms of intraseasonal variability in ocean dynamic sea level (ζ) in the Persian Gulf over 2002–2015. Empirical orthogonal function (EOF) analysis applied to altimetry data reveals a basin-wide, single-signed intraseasonal fluctuation that contributes importantly to ζ variance in the Persian Gulf at monthly to decadal timescales. An EOF analysis of Gravity Recovery and Climate Experiment (GRACE) observations over the same period returns a similar large-scale mode of intraseasonal variability, suggesting that the basin-wide intraseasonal ζ variation has a predominantly barotropic nature. A linear barotropic theory is developed to interpret the data. The theory represents Persian-Gulf-average ζ ($\bar{\zeta}$) in terms of local freshwater flux, barometric pressure, and wind stress forcing, as well as ζ at the boundary in the Gulf of Oman. The theory is tested using a multiple linear regression with these freshwater flux, barometric pressure, wind stress, and boundary ζ quantities as input, and $\bar{\zeta}$ as output. The regression explains $70\% \pm 9\%$ (95% confidence interval) of the intraseasonal $\bar{\zeta}$ variance. Numerical values of regression coefficients computed empirically from the data are consistent with theoretical expectations from the theory. Results point to a substantial non-isostatic response to surface loading. The Gulf of Oman ζ boundary condition shows lagged correlation with ζ upstream along the Indian Subcontinent, Maritime Continent, and equatorial Indian Ocean, suggesting a large-scale Indian-Ocean influence on intraseasonal $\bar{\zeta}$ variation mediated by coastal and equatorial waves, and hinting at potential predictability. This study highlights the value of GRACE for understanding sea level in an understudied marginal sea.

28 1. Introduction

29 The Persian Gulf¹ is a semi-enclosed marginal sea of the Indian Ocean (Figure 1). It connects to
30 the Arabian Sea to the southeast through the Strait of Hormuz and the Gulf of Oman. The Persian
31 Gulf is shallow and broad, with an average depth of ~ 30 m and a surface area of $\sim 2.2 \times 10^5$ km².
32 It is subject to an arid, subtropical climate, and is bounded to the southwest by the Arabian Desert
33 and by the Zagros mountains to the northeast.

34 Past studies establish the basic physical oceanography of the Persian Gulf using data and models
35 (Chao et al., 1992; Emery, 1956; Johns et al., 1999, 2003; Kämpf and Sadrinasab, 2006; Reynolds,
36 1993; Thoppil and Hogan, 2010; Swift and Bower, 2003; Yao and Johns, 2010). We outline some
37 of the salient features for context. The region is forced year-round by north-northwesterly surface
38 winds ('shamal', speeds 3–6 m s⁻¹). Evaporation (~ 2 m y⁻¹) far exceeds precipitation and runoff
39 (~ 0.2 m y⁻¹), resulting in an inverse-estuarine circulation—fresher, warmer buoyant waters inflow
40 near the surface through the Strait of Hormuz largely along the coast of Iran, whereas saltier, colder,
41 denser waters outflow near the bottom mainly along the coast of the United Arab Emirates. The
42 basin-scale circulation is demarcated by a thermal front across the Persian Gulf between Qatar and
43 Iran. Northwest of the front, there is equatorward flow along Saudi Arabia driven by wind-forced
44 downwelling at the coast and buoyant river discharge from the Tigris, Euphrates, and other rivers
45 at the head of the Persian Gulf. To the southeast, there exists a large-scale counterclockwise
46 circulation, maintained by exchanges through the Strait of Hormuz, and evaporation, cooling, and
47 sinking of water masses in shallow regions along the southern Persian Gulf. Mesoscale eddies are
48 common, especially during boreal summer, when they are shed from the Iranian Coastal Jet due to
49 baroclinic instability. For more details, interested readers are directed to the papers cited above.

¹The name of this body of water is subject to dispute. It is also known as the Arabian Gulf or the Gulf. We use the name Persian Gulf following the conventions of the International Hydrographic Organization and the United Nations.

50 The Persian Gulf is one of the world ocean's busiest waterways, due to its vast oil and gas stores,
51 which are of longstanding geopolitical, economic, and military interest (al-Chalabi, 2007; Barnes
52 and Myers Jaffe, 2006; Larson, 2007). Bordering eight nations, the Persian Gulf is also home to
53 large coastal populations and major coastal cities including Dubai, Abu Dhabi, and Doha, which
54 are exposed to risk of flooding and inundation related to sea-level change (Al-Jeneid et al., 2008;
55 Lafta et al., 2020). Kopp et al. (2014, 2017) project that mean sea level will rise by 44–108 cm
56 between 2000 and 2100 in Bahrain under the Representative Concentration Pathway 8.5 forcing
57 scenario (66% confidence). This would threaten $\sim 10\text{--}15\%$ ($\sim 80\text{--}100 \text{ km}^2$) of Bahrain's surface
58 area (Al-Jeneid et al., 2008). Such numbers emphasize the importance of understanding sea-level
59 changes in the Persian Gulf. However, projections of mean sea-level rise on multidecadal and longer
60 timescales (Kopp et al., 2014, 2017) alone are insufficient to anticipate future coastal flood risk.
61 Also important are sea-level fluctuations at decadal and shorter periods, which can superimpose
62 on longer-term changes, temporarily ameliorating or exacerbating coastal risk (Burgos et al., 2018;
63 Dangendorf et al., 2016; Long et al., 2020; Ray and Foster, 2016; Sweet et al., 2017). This
64 motivates a detailed investigation of mean sea-level variation in the Persian Gulf on decadal and
65 shorter timescales—what are the dominant magnitudes, scales, and mechanisms?

66 Past studies on Persian Gulf mean sea level largely focus on seasonal cycles and decadal trends
67 (Al-Subhi, 2010; Alothman et al., 2014; Ayhan, 2020; Barzandeh et al., 2018; El-Gindy, 1991;
68 El-Gindy and Eid, 1997; Hassanzadeh et al., 2007; Hosseinibalam et al., 2007; Sharaf El Din,
69 1990; Siddig et al., 2019; Sultan et al., 1995a, 2000). Sultan et al. (1995a) consider monthly
70 relative sea level during 1980–1990 from two tide gauges on the Saudi Arabia coast. They find
71 that 80% of the overall monthly data variance is explained by the seasonal cycle, which has an
72 amplitude of $\sim 10 \text{ cm}$ and peaks in boreal summer. These authors argue that 75% of the seasonal
73 variance in sea level reflects an inverted-barometer response to a $\sim 10\text{-mb}$ -amplitude seasonal cycle

74 in local surface air pressure, and that the remaining 25% of seasonal variance represents steric
 75 variability owing to density fluctuations. Other studies targeting different regions, tide gauges, and
 76 time periods confirm this basic result that inverted-barometer and steric effects make primary and
 77 secondary contributions, respectively, to the large-scale seasonal cycle in Persian Gulf sea level,
 78 but also suggest that local wind effects are important in some places (Al-Subhi, 2010; Barzandeh
 79 et al., 2018; El-Gindy, 1991; El-Gindy and Eid, 1997; Hassanzadeh et al., 2007; Hosseinibalam et
 80 al., 2007; Sharaf El Din, 1990; Sultan et al., 2000). Alothman et al. (2014) interrogate monthly
 81 relative sea level over 1979–2007 based on 15 tide-gauge records from Bahrain, Saudi Arabia, and
 82 Iran, along with measurements of vertical land motion from 6 Global Positioning System (GPS)
 83 stations in Bahrain, Saudi Arabia, and Kuwait. They determine that regional relative sea level rose
 84 by $2.2 \pm 0.5 \text{ mm y}^{-1}$ over that time. These authors find that one-third of the increase ($0.7 \pm 0.6 \text{ mm}$
 85 y^{-1}) was due to crustal subsidence, possibly related to groundwater pumping and oil extraction
 86 (Amin and Bankher, 1997), and the remaining two-thirds ($1.5 \pm 0.8 \text{ mm y}^{-1}$) was due to geocentric
 87 sea-level changes. Sultan et al. (2000) calculate a more muted relative sea-level trend (1.7 mm
 88 y^{-1}) based on 9 tide-gauge records from Saudi Arabia over 1980–1994, while Siddig et al. (2019)
 89 estimate a larger geocentric sea-level trend ($3.6 \pm 0.4 \text{ mm y}^{-1}$) from altimetry data averaged over
 90 the Persian Gulf during 1993–2018, consistent with reports of a global sea-level acceleration in
 91 recent decades (Nerem et al., 2018; Dangendorf et al., 2019; Frederikse et al., 2020).

92 Omitted from past works on Persian Gulf mean sea level is exploration of nonseasonal sea-level
 93 variation. This is an important omission, since nonseasonal variations in general, and in particular
 94 intraseasonal variations, contribute importantly to mean sea-level variance over the Persian Gulf on
 95 monthly to decadal timescales. For example, consider the time series of monthly ocean dynamic
 96 sea level from satellite-altimetry data averaged over the Persian Gulf during 2002–2015 shown in
 97 Figure 2. Filters are applied to the data to emphasize variability on different timescales, and global-

mean sea level and the inverted-barometer effect are removed. Nonseasonal fluctuations explain 52% of the monthly data variance, and intraseasonal fluctuations (with ~ 2 –6-month periods) alone account for 46% of the overall data variance. The altimetric time series of intraseasonal sea level averaged over the Persian Gulf also explains 51% of the intraseasonal variance in relative sea level averaged across 5 tide gauges from Iran and Bahrain during the overlapping period 2002–2006 (Figure 2). This exploratory analysis suggests that large-scale intraseasonal fluctuations make important contributions to ocean dynamic sea-level variance across the Persian Gulf during the altimeter era, motivating a more in-depth investigation.

Here we investigate the magnitudes, scales, and mechanisms of intraseasonal sea-level variability in the Persian Gulf through an analysis of satellite observations and other data. The remainder of the paper is structured as follows: in section 2, we describe the data; in section 3, we establish the horizontal scales and vertical structure of the dominant intraseasonal sea-level variation in the Persian Gulf; in section 4, we use dynamical theory, linear regression, and correlation analysis to identify the main local and nonlocal forcing mechanisms and ocean dynamics responsible for driving intraseasonal variations in Persian Gulf sea level and their relation to large-scale circulation and climate in the Indian Ocean; we conclude with a summary and discussion in section 5.

2. Materials and Methods

a. Ocean dynamic sea level from satellite altimetry

We use version 2.0 of the sea-level essential climate variable product from the European Space Agency Climate Change Initiative (Legeais et al., 2018; Quartly et al., 2017). Data were downloaded from the Centre for Environmental Data Analysis on 18 April 2020. (All data sources are indicated in Table 1.) The multi-satellite merged geocentric sea-level anomalies are given on a

0.25° global spatial grid and a monthly time increment during 1993–2015. These data extend and update the earlier version 1.1 product (Ablain et al., 2015). The dynamic atmospheric correction is applied, which involves removing the ocean’s dynamic barotropic response to wind and pressure forcing at shorter periods < 20 days and its isostatic response to pressure forcing at longer periods > 20 days from the data (Carrère and Lyard, 2003; Carrère et al. 2016). (The dynamic ocean response to these forcings at the periods of interest to this study are retained in the data.) For more details on the geophysical corrections, orbit solutions, altimeter standards, and error budgets, see Quartly et al. (2017) and Legeais et al. (2018). We remove the time series of global-mean geocentric sea-level values from every grid cell. Assuming that gravitational, rotational, and deformational effects are negligible (Gregory et al., 2019), the resulting sea-level anomalies represent ocean dynamic sea-level anomalies². We use the data from May 2002 to September 2015, since this is the period of overlap between this altimeter data set and the Gravity Recovery and Climate Experiment (GRACE), which is used for interpretation and described below. Following Gregory et al. (2019), we use ζ to denote ocean dynamic sea level.

This paper focuses on intraseasonal variability. To isolate intraseasonal behavior, we process the data as follows. We use least squares to estimate the seasonal cycle (annual and semi-annual sinusoids) and linear trend in the data over the study period. We then remove these seasonal and trend contributions from the original data to create a time series of nonseasonal residuals. Next, we apply a Gaussian smoother with a 3-month half window to these nonseasonal residuals. Finally, we subtract this low-pass-filtered time series from the nonseasonal residuals to create a record of intraseasonal fluctuations, which is the object of our study. We delete the first and last 6 months of the intraseasonal time series to avoid edge effects. This filter passes $> 90\%$ of the power at periods

²Ocean dynamic sea level refers to the local height of the sea surface above the geoid with the inverted-barometer correction applied (Gregory et al., 2019).

142 $\lesssim 8$ months and stops $> 70\%$ of the power at periods $\gtrsim 15$ months. See Figure 2 for an example of
143 this filtering applied to altimetry averaged over the Persian Gulf.

144 *b. Manometric sea level from satellite gravimetry*

145 We consider data from GRACE and GRACE Follow-On (Landerer et al., 2020; Watkins et al.,
146 2015; Wiese et al., 2016). Mass grids were downloaded from the National Aeronautics and Space
147 Administration Jet Propulsion Laboratory on 15 April 2020 (data version JPL RL06M.MSCNv02).
148 The data are processed using 3° spherical-cap mass-concentration blocks for the gravity-field basis
149 functions. For more details on the estimation process, spatial constraints, scale factors, and leakage
150 errors, see Watkins et al. (2015). The data are defined on a 0.5° global spatial grid, but the satellite
151 measurement do not resolve processes with spatial scales $\lesssim 300$ km. We use the version of the
152 data with the coastline resolution improvement filter applied (Wiese et al., 2016). The grids are
153 defined at irregular, quasi-monthly increments, and have gaps. For example, battery management
154 issues caused multi-month data gaps in the final years of GRACE, and there is a ~ 1 -y data gap
155 between the end of GRACE coverage and the beginning of the GRACE Follow-On record. We
156 linearly interpolate the available ocean mass grids onto regular monthly increments from May 2002
157 through September 2015. The data have units of equivalent water thickness. After correcting for
158 global air-pressure effects, these data reflect manometric sea-level anomalies³. To isolate dynamic
159 manometric sea-level anomalies associated with internal ocean mass redistribution, we subtract the
160 time series of barystatic sea level⁴ from the data at every oceanic grid cell. Intraseasonal variations
161 are isolated through filtering methods described earlier. Following Gregory et al. (2019), we use
162 R_m to indicate manometric sea level, with its dynamic nature understood.

³Manometric sea-level changes indicate sea-level changes due to changes in the local mass of the ocean per unit area (Gregory et al., 2019).

⁴Barystatic sea-level changes refer to global-mean sea-level changes due to net addition or subtraction of water mass to or from the global ocean (Gregory et al., 2019).

c. Relative sea level from tide gauges

We also use monthly mean relative sea level⁵ from tide-gauge records in the Persian Gulf that overlap with our study period (Table 2). Data were downloaded from the Permanent Service for Mean Sea Level database on 1 July 2019 (PSMSL, 2019; Holgate et al., 2013). The data from Mina Sulman in Manama, Bahrain represent the only record from the Persian Gulf in the PSMSL database with a complete benchmark datum history (so-called revised local reference data). To consider large-scale regional behavior, we also study a careful selection of records without continuous datum histories (so-called metric data). Namely, we use the data from Emam Hassan, Bushehr, Kangan, and Shahid Rajaee in Iran⁶. We consider the data over 2002–2006, since earlier times predate our study, and later times feature no tide-gauge data (Table 2). The data from Emam Hassan before November 2002 are omitted due to a data gap that coincided with an apparent datum shift (Allothman et al., 2014). We adjust each record for the inverted-barometer effect using reanalysis surface air pressure (see below). Next, we remove the seasonal cycle and linear trend from each adjusted time series. We then average together these nonseasonal time series to create a regional composite of adjusted relative sea level. Finally, we isolate intraseasonal variability by computing and then removing a low-pass-filtered version of the regional composite. The resulting time series is shown in Figure 2. To the extent that global-mean sea-level changes and gravitational, rotational, and deformational effects are unimportant on these scales, this composite time series represents tide-gauge-based intraseasonal regional ζ variability.

⁵Relative sea level is the height of the sea surface relative to the solid Earth (Gregory et al., 2019).

⁶Metric data from other Persian Gulf locations are also available in the PSMSL database. However, we determined that these records were unsuitable for our analysis. Five records from the United Arab Emirates, Qatar, and Iraq are short and predate our study period. A dozen records from Saudi Arabia were operated by the Saudi Arabian Oil Company and situated on oil platforms, and are therefore potentially unstable.

d. Surface forcing

We use gridded observations, atmospheric reanalyses, and flux estimates to interpret the data from altimetry, GRACE, and tide gauges. For all fields, we compute intraseasonal anomalies during 2002–2015 from the available monthly values, as with the altimetry and GRACE.

We use monthly wind stress and barometric pressure from the European Centre for Medium Range Weather Forecasts Reanalysis Interim (ERA-Interim; Dee et al., 2011). Fields were downloaded from the Woods Hole Oceanographic Institution (WHOI) Community Storage Server on 7 January 2019. Values are defined on a 0.75° global spatial grid from January 1979 to October 2018.

We use monthly evaporation from version 3 of the the Objectively Analyzed air-sea Fluxes project (OAFlux; Yu and Weller, 2007). Fields were downloaded from WHOI servers on 13 November 2019. Values are defined on a 1° global spatial grid from January 1958 to December 2018.

We use monthly precipitation from version 2.3 of the Global Precipitation Climatology Project (GPCP; Adler et al., 2003). Fields were downloaded from National Oceanic and Atmospheric Administration Earth System Research Laboratory and Physical Sciences Laboratory on 16 April 2020. Values are defined on a 2.5° global spatial grid from January 1979 to the present.

We use monthly river runoff from the Japanese 55-year atmospheric reanalysis surface data set for driving ocean–sea-ice models (JRA55-do; Tsujino et al., 2018). Fields were downloaded from servers at the Hokkaido University Graduate School of Environmental Science on 21 August 2020. Values are defined on a 0.25° global coastal grid from January 1958 to December 2017.

3. Horizontal scales and vertical structure of ζ variability

Past studies use satellite altimetry and tide gauges to study seasonal cycles and decadal trends in the Persian Gulf (Al-Subhi, 2010; Alothman et al., 2014; Ayhan, 2020; El-Gindy, 1991; El-Gindy and Eid, 1997; Hassanzadeh et al., 2007; Hosseinibalam et al., 2007; Sharaf El Din, 1990; Siddig

et al., 2019; Sultan et al., 1995a, 2000). Here we examine intraseasonal variability in the Persian Gulf using satellite data, including altimetry but also gravimetry, and tide gauges.

We motivated this study with an exploratory data analysis earlier in the introduction. We found that roughly half of the monthly ζ variance from altimetry averaged over the Persian Gulf during 2002–2015 was concentrated at intraseasonal periods, and that the Persian-Gulf-average altimetric time series of intraseasonal ζ ($\bar{\zeta}$) explained about half of the variance in a composite time series of intraseasonal ζ from coastal tide gauges (Figure 2). These results show that intraseasonal fluctuations contribute importantly to large-scale ζ variability over the Persian Gulf at monthly to decadal periods, and that intraseasonal fluctuations measured locally at the coast largely reflect spatially coherent, basin-wide behavior.

To explore intraseasonal ζ in more detail, we apply empirical orthogonal function (EOF) analysis to altimetry data over the Persian Gulf. We identify the spatial structures and temporal behaviors of the orthogonal modes of intraseasonal variability by solving for the eigenvalues and eigenvectors of the covariance matrix of the altimetry data over the Persian Gulf. The eigenvectors correspond to the spatial structures and the eigenvalues indicate the amounts of data variance explained by the various modes. The temporal behaviors of the modes are described by principal-component time series, which are determined by projecting the respective eigenvectors onto the data (von Storch and Zwiers, 1999).

The leading mode, which explains 52% of the intraseasonal data variance over the Persian Gulf, is summarized in Figures 3 and 4. It shows a single-signed spatial structure (Figure 3a), indicating basin-wide variation and wholesale raising and lowering of ζ over the Persian Gulf. This is consistent with our earlier finding that the $\bar{\zeta}$ time series from altimetry explains 51% of the variance in the regional composite from tide gauges at intraseasonal timescales (Figure 2). Indeed, this mode's principal-component time series (Figure 4) is perfectly correlated with the $\bar{\zeta}$

time series from altimetry (correlation coefficient > 0.99). The leading mode from a complex-valued (Hilbert) EOF analysis explains the same amount of data variance (not shown). This means that out-of-phase relationships between ζ in different parts of the Persian Gulf related to signal propagation are unimportant to this mode, and that this dominant ζ variation reflects an in-phase standing mode of oscillation across the region on these timescales.

The spatial structure is also nonuniform (Figure 3a). Magnitudes increase from southeast to northwest across the region, with smaller values (1–3 cm) observed along the United Arab Emirates, Qatar, Bahrain, and southern Iran, and larger values (3–5 cm) apparent off Saudi Arabia, Kuwait, Iraq, and northern Iran. This basin-scale gradient could reflect wind setup related to strengthening or weakening of the region’s prevailing north-northwesterlies. The strongest amplitudes (> 5 cm) are detected off Kuwait and Iraq, near the mouths of the Tigris, Euphrates, and Karun rivers. Values in this region are highest at the coast and decay offshore, possibly indicating trapped signals driven by buoyant river discharge. There is also spatial structure in the amount of local data variance explained by this mode: whereas 50–80% of local ζ data variance is explained over the interior in the northwestern Persian Gulf, $< 30\%$ is explained in the southwest off Qatar, Bahrain, and the United Arab Emirates (Figure 3b). This suggests important local-scale ζ variability along the southwest coast that is unrelated to the broader-scale behavior resolved by this mode.

The ζ response to surface forcing is often described in terms of barotropic (depth-independent) and baroclinic (depth-dependent) adjustments (e.g., Vinogradova et al., 2007). Given the latitude of the Persian Gulf, and the spatiotemporal scales under investigation, basic scaling arguments (Gill and Niiler, 1973; Piecuch et al., 2019) suggest that this mode of ζ variation should be essentially barotropic in nature. For a purely barotropic ocean response, changes in sea level (or subsurface pressure) are mirrored by changes in ocean bottom pressure (Bingham and Hughes, 2008; Vinogradova et al., 2007). Hence, if the leading mode of ζ variability from altimetry

(Figure 3, 4) reflects a predominantly barotropic response, then similar R_m variability should be apparent in GRACE.

To test this hypothesis, we apply EOF analysis to the GRACE R_m grids over the Persian Gulf. The results are shown in Figures 4 and 5. The leading mode, which explains 88% of the intraseasonal GRACE data variance in the Persian Gulf, shows a single-signed spatial pattern, such that variability increases from 1–2 cm in the southeastern Persian Gulf to 3–4 cm in the northwest (Figure 5a). Relatively more local R_m data variance is explained ($> 80\%$) to the north and west, while comparatively less is explained (50–70%) in the southeast (Figure 5b). These patterns from GRACE are qualitatively similar to those from altimetry, but there are quantitative differences (cf. Figures 3, 5). For example, the mode from altimetry exhibits larger amplitudes and richer, more detailed spatial structures than the mode from GRACE (Figures 3a, 5a), whereas the leading GRACE mode explains relatively more data variance compared to the leading altimetry mode (Figures 3b, 5b). These discrepancies probably partly reflect the coarser resolution (and reduced effective spatial degrees of freedom) of GRACE, but could also indicate baroclinic processes or data errors (e.g., residual leakage of terrestrial signals into the GRACE ocean grids).

Such differences notwithstanding, results in Figures 3 and 5 suggest that GRACE and altimetry capture facets of the same underlying mode of intraseasonal variation. This suggestion is corroborated by the principal components of the leading EOF modes determined from GRACE and altimetry, which are highly correlated (correlation coefficient of ~ 0.7 ; Figure 4). We also apply maximum covariance analysis (MCA) jointly to altimetry ζ and GRACE R_m data, whereby the eigenvalues and eigenvectors of the cross-covariance matrix between the two data sets are determined (von Storch and Zwiers, 1999). The leading eigenvectors and principal components determined jointly through MCA are identical to those determined separately through EOF analysis, and the gravest MCA mode explains $> 99\%$ of the joint covariance between altimetry and

GRACE data (not shown). This suggests that the leading modes of regional ζ and R_m variation are coupled to one another, and reflect a dominant barotropic response.

4. Forcing mechanisms and ocean dynamics

In the previous section, we established a basin-wide barotropic variation of the Persian Gulf on intraseasonal timescales. Here we use analytical theory, linear regression, and correlation analysis to identify the forcing and dynamics responsible for this mode.

a. Linear barotropic model

The leading mode of intraseasonal variability identified previously exhibits higher-order spatial structure (Figures 3, 5). However, the lowest-order spatial feature is that of a horizontally uniform fluctuation. For example, the time series of intraseasonal $\bar{\zeta}$ from altimetry explains 93% of the variance associated with the first altimetric EOF mode (Figures 2–4). Thus, we formulate a linear model for a horizontally uniform barotropic variation of the Persian Gulf. Our formulation largely follows Volkov et al. (2016), who use a similar model to consider ζ in the Black Sea. The equations for conservation of volume within the Persian Gulf and conservation of momentum along the Strait of Hormuz are

$$S\bar{\zeta}_t = S\bar{q} + \frac{S}{\rho g}\bar{p}_t + vWH, \quad (1)$$

$$v_t = -g\zeta_y + \frac{1}{\rho H}\tau - \frac{r}{H}v. \quad (2)$$

Here S is surface area of the Persian Gulf, overbar is spatial average over the Persian Gulf, q is precipitation plus runoff minus evaporation, p is barometric pressure, v is average velocity along the Strait of Hormuz into the Persian Gulf (positive values increase the volume of the Persian Gulf), W and H are the width and depth of the Strait of Hormuz, respectively, τ is wind stress along the Strait of Hormuz (positive in the direction of the Persian Gulf), r is a constant friction coefficient,

298 g is gravity, ρ is seawater density, and subscripts t and y denote partial differentiation in time and
 299 the along-strait direction, respectively. Note that, since we express Eqs. (1) and (2) in terms of ζ ,
 300 forcing by p appears in the continuity equation rather than in the momentum equation, and takes
 301 on an analogous form to the q forcing, such that, as noted by Gill (1982), forcing by a depression
 302 of 10 mb would be canceled out by 10 cm of precipitation (cf. also Ponte, 2006). All symbols are
 303 described in Table 3 and representative values are given when appropriate.

304 We assume ζ , v , q , p , and τ take wave solutions of the form $\exp(-i\omega t)$ with angular frequency
 305 ω and $i \doteq \sqrt{-1}$. Integrating the momentum equation over the length L of the Strait of Hormuz, and
 306 rearranging to solve for $\bar{\zeta}$ gives

$$\bar{\zeta} = \left[\zeta_0 + \frac{L}{\rho g H} \tau + \frac{(\lambda - i\omega)}{\sigma^2} \bar{q} - i\omega \frac{(\lambda - i\omega)}{\sigma^2} \frac{\bar{p}}{\rho g} \right] \left/ \left[1 - \frac{\omega^2}{\sigma^2} - i \frac{\lambda \omega}{\sigma^2} \right] \right., \quad (3)$$

307 where ζ_0 represents ζ at the boundary outside the Strait of Hormuz in the Gulf of Oman, and
 308 we define $\sigma^2 \doteq WHg/SL$ and $\lambda \doteq r/H$. Physically, $1/\lambda$ is a friction timescale and $1/\sigma$ is a
 309 Helmholtz resonance timescale determined by the shape of the Persian Gulf and Strait of Hormuz.
 310 (We determine that $1/\sigma \approx 15$ hours, which is small compared to the intraseasonal timescales of
 311 interest, so we do not expect a resonant response.) Equivalently, we can write Eq. (3) in the polar
 312 complex plane as

$$\bar{\zeta} = z_{\zeta_0} \exp(i\theta_{\zeta_0}) \zeta_0 + z_{\tau} \exp(i\theta_{\tau}) \tau + z_{\bar{q}} \exp(i\theta_{\bar{q}}) \bar{q} + z_{\bar{p}} \exp(i\theta_{\bar{p}}) \bar{p}, \quad (4)$$

313 where

$$\theta_{\zeta_0} \doteq \arctan\left(\frac{\lambda \omega}{\sigma^2 - \omega^2}\right), \quad (5)$$

$$z_{\zeta_0} \doteq \left[\left(1 - \frac{\omega^2}{\sigma^2}\right)^2 + \left(\frac{\lambda \omega}{\sigma^2}\right)^2 \right]^{-1/2}, \quad (6)$$

$$\theta_{\tau} \doteq \arctan\left(\frac{\lambda \omega}{\sigma^2 - \omega^2}\right), \quad (7)$$

$$z_\tau \doteq \left[\left(1 - \frac{\omega^2}{\sigma^2} \right)^2 + \left(\frac{\lambda\omega}{\sigma^2} \right)^2 \right]^{-1/2} \left(\frac{L}{\rho g H} \right), \quad (8)$$

$$\theta_{\bar{q}} \doteq \arctan \left(\frac{\lambda\omega}{\sigma^2} - \frac{\omega}{\lambda} + \frac{\omega^3}{\sigma^2\lambda} \right), \quad (9)$$

$$z_{\bar{q}} \doteq \frac{\lambda}{\sigma^2} \left[1 + \left(\frac{\lambda\omega}{\sigma^2} - \frac{\omega}{\lambda} + \frac{\omega^3}{\sigma^2\lambda} \right)^2 \right]^{1/2} \left[\left(1 - \frac{\omega^2}{\sigma^2} \right)^2 + \left(\frac{\lambda\omega}{\sigma^2} \right)^2 \right]^{-1}, \quad (10)$$

$$\theta_{\bar{p}} \doteq \arctan \left[\left(\frac{\omega}{\lambda} - \frac{\omega^3}{\lambda\sigma^2} - \frac{\omega\lambda}{\sigma^2} \right)^{-1} \right], \quad (11)$$

$$z_{\bar{p}} \doteq \frac{1}{\rho g} \frac{\lambda\omega}{\sigma^2} \left[1 + \left(\frac{\omega}{\lambda} - \frac{\lambda\omega}{\sigma^2} - \frac{\omega^3}{\sigma^2\lambda} \right)^2 \right]^{1/2} \left[\left(1 - \frac{\omega^2}{\sigma^2} \right)^2 + \left(\frac{\omega\lambda}{\sigma^2} \right)^2 \right]^{-1}, \quad (12)$$

In other words, according to Eq. (4), $\bar{\zeta}$ is a linear superposition of the ζ_0 , τ , \bar{q} , and \bar{p} forcing terms, each scaled by an amount z_j and rotated through a phase θ_j , where $j \in \{\zeta_0, \tau, \bar{q}, \bar{p}\}$. We estimate theoretical values for the scaling factors z_j and phase angles θ_j by averaging Eqs. (5)–(12) over the ω range from $2\pi / (6 \text{ months})$ to $2\pi / (2 \text{ months})$ using numerical values for the scalar coefficients λ , σ , L , ρ , g , and H from Table 3. These theoretical values are tabulated in Table 4.

b. Multiple linear regression analysis

To test whether the model described by Eqs. (1)–(12) is informative for understanding observed intraseasonal $\bar{\zeta}$ variability, we perform a multiple linear regression. We model $\bar{\zeta}$ from altimetry as

$$\bar{\zeta} = a_{\zeta_0}\zeta_0 + b_{\zeta_0}\mathcal{H}(\zeta_0) + a_\tau\tau + b_\tau\mathcal{H}(\tau) + a_{\bar{q}}\bar{q} + b_{\bar{q}}\mathcal{H}(\bar{q}) + a_{\bar{p}}\bar{p} + b_{\bar{p}}\mathcal{H}(\bar{p}) + \varepsilon, \quad (13)$$

where \mathcal{H} is the Hilbert transform, the a_j and b_j are real constants, and ε is the residual. We include Hilbert transforms of the various forcings in the regression to allow for possible phase lags between the forcing and the response, as indicated by Eq. (4). We estimate the z_j and θ_j from Eq. (4) from the a_j and b_j in Eq. (13) using properties of Hilbert transforms and trigonometric identities as

$$\theta_j = \arctan(b_j/a_j), \quad (14)$$

$$z_j = \sqrt{a_j^2 + b_j^2}. \quad (15)$$

334 We evaluate Eq. (13) using least squares. For ζ_0 , we use ζ from altimetry averaged over
 335 shallow regions (< 200 m) of the northern Gulf of Oman outside the Strait of Hormuz ($57\text{--}60^\circ\text{E}$,
 336 $25\text{--}28^\circ\text{N}$). For τ , we use along-strait wind stress (315°T) from ERA-Interim averaged over the
 337 Strait of Hormuz ($54\text{--}57.8^\circ\text{E}$, $22.9\text{--}27.4^\circ\text{N}$). For \bar{q} , we use precipitation from GPCP plus river
 338 runoff from JRA55-do minus evaporation from OAFlux averaged over the Persian Gulf ($45\text{--}55^\circ\text{E}$,
 339 $24\text{--}32^\circ\text{N}$). For \bar{p} , we use barometric pressure from ERA-Interim averaged over the Persian Gulf
 340 ($48\text{--}54.8^\circ\text{E}$, $24.4\text{--}29.6^\circ\text{N}$). Uncertainties are estimated using 10 000 iterations of bootstrapping
 341 (Efron and Hastie, 2016).

342 Results of the multiple linear regression are summarized in Figure 6. The regression model
 343 [(13)] explains $70\% \pm 9\%$ (95% confidence interval) of the variance in the $\bar{\zeta}$ data (Figure 6a). This
 344 suggests that Eqs. (1) and (2) represent the dominant physics, and that $\bar{\zeta}$ variability can be largely
 345 understood in terms of local surface forcing by τ , \bar{q} , and \bar{p} and nonlocal boundary forcing by ζ_0 .
 346 In Figure 6b, we break down the relative contributions of the different forcing terms. The primary
 347 driver of $\bar{\zeta}$ is nonlocal forcing by ζ_0 , which explains $50\% \pm 12\%$ of the $\bar{\zeta}$ variance. Local forcing
 348 by τ , \bar{q} , and \bar{p} plays a secondary role. Individually, τ explains $16\% \pm 9\%$, \bar{q} explains $5\% \pm 9\%$,
 349 and \bar{p} explains $10\% \pm 8\%$ of the $\bar{\zeta}$ variance. Surface loading (the combination of \bar{q} and \bar{p} forcing)
 350 explains $14\% \pm 11\%$ of the variance in the data. Collectively, all three local forcing factors taken
 351 together account for $27\% \pm 14\%$ of the $\bar{\zeta}$ variance.⁷

⁷The variance contributions of the individual predictors are not entirely additive, since they are not wholly independent and there is some correlation between them. However, the relative roles of the respective forcings can nevertheless be meaningfully estimated (albeit with uncertainty) because the least-squares problem is generally well posed. After normalizing the predictors to unit variance, the condition number of their covariance matrix is 3.3. This is on the same order as the range of 1.4–2.5 (99% confidence interval) we determine through repeated simulations of four independent random, standard-normal time series (and their Hilbert transforms) with the same length as the observations (not shown).

Regression coefficients computed empirically from the data are consistent with values expected theoretically from first principles (Table 4). For example, the linear regression yields a scaling factor of $1.5 \pm 0.5 \text{ m Pa}^{-1}$ and a phase angle of 30 ± 25 degrees between τ and $\bar{\zeta}$. This is consistent with the theoretical ranges of $1.0\text{--}1.3 \text{ m Pa}^{-1}$ and $5\text{--}38$ degrees anticipated from Eqs. (7) and (8). The regression analysis also suggests a substantial departure from the inverted-barometer response, manifested in a scaling of $0.8 \pm 0.5 \text{ cm mb}^{-1}$ and a phase of 65 ± 52 degrees between \bar{p} and $\bar{\zeta}$. This overlaps with the ranges of $0.1\text{--}0.5 \text{ cm mb}^{-1}$ and $56\text{--}87$ degrees expected from Eqs. (11) and (12). (Recall that the altimeter data have been adjusted for an inverted barometer and that our theory was developed for ζ , which has the inverted-barometer effect already removed.) This provides evidence that the results of the multiple linear regression indicate true causal relationships between forcing and response.

Regression results and analytical theory suggest that these relationships can be out of phase, such that the forcings lead the response by a significant amount (Table 4). To quantify the importance of out-of-phase behavior, we perform another multiple linear regression analysis, this time omitting Hilbert transforms and forcing by p from the input [cf. Eq. (13)]. Physically, this alternative regression model assumes an equilibrium response, and corresponds to the steady state ($\omega \rightarrow 0$) limit of the governing equations, viz. [cf. Eq. (3)],

$$\bar{\zeta} = \zeta_0 + \frac{L}{\rho gh} \tau + \frac{\lambda}{\sigma^2} \bar{q}. \quad (16)$$

This alternate model accounts for slightly less of the $\bar{\zeta}$ data variance ($62\% \pm 10\%$; 95% confidence interval). This result demonstrates that a majority of the $\bar{\zeta}$ data variance explained by the original multiple linear regression model [Eq. (13)] is attributable to equilibrium processes and in-phase (or antiphase) relationships between the forcing and the response, but also that allowing for transient processes [the time derivatives in Eqs. (1) and (2)] and more general phase relationships between

forcing and response leads to a modest, but significant, improvement in terms of explaining $\bar{\zeta}$ data variance.

c. Relation to Indian Ocean circulation and climate, and potential predictability

Nonlocal forcing by ζ_0 is the most important contributor to $\bar{\zeta}$ variability (Figure 6b). What is the nature of these fluctuations at the boundary in the Gulf of Oman? How do they relate to larger-scale circulation and climate? To clarify their origin, we compute correlation coefficients between ζ_0 and either ζ or its Hilbert transform $\mathcal{H}(\zeta)$ at every altimetric grid point over the Indian Ocean. Correlations between ζ_0 and ζ identify regions where ζ is in phase or anti-phase (i.e., 180 degrees out of phase) with ζ_0 , whereas correlations between ζ_0 and $\mathcal{H}(\zeta)$ indicate regions where ζ is in quadrature (90 degrees out of phase) or anti-quadrature (270 degrees out of phase) with ζ_0 .

In general, ζ_0 is uncorrelated with ζ and $\mathcal{H}(\zeta)$ away from the coast and the equator (Figure 7), suggesting that ζ_0 is unrelated to the dominant ζ variability in these open-ocean regions. However, we observe patterns of significant correlation and anti-correlation along the coast and equator. For example, ζ_0 is correlated with ζ along Pakistan, western India, and Sri Lanka; correlated with $\mathcal{H}(\zeta)$ along eastern India, Bangladesh, and Myanmar; correlated with $\mathcal{H}(\zeta)$ and anti-correlated with ζ along Thailand, Malaysia, and Sumatra; and anti-correlated with $\mathcal{H}(\zeta)$ along the western equatorial Indian Ocean between Somalia and the Maldives (Figure 7).

These patterns suggest wave propagation along equatorial and coastal waveguides. For example, the correlation between ζ_0 and $\mathcal{H}(\zeta)$ along Bangladesh suggests that ζ_0 lags ζ in this region by 90 degrees (one quarter of a period), whereas anti-correlation between ζ_0 and $\mathcal{H}(\zeta)$ in the western equatorial Indian Ocean hints that regional ζ leads ζ_0 by 270 degrees (three quarters of a period). Supposing propagation is eastward along the equator and counterclockwise along the coast (in the Northern Hemisphere), and assuming intraseasonal periods of 60–180 days, we estimate that these

phase leads and lags imply propagation speeds of $\sim 1\text{--}3 \text{ m s}^{-1}$. These values are consistent with basic expectations for equatorial waves and coastally trapped waves (e.g., Gill, 1982; Hughes et al., 2019). Indeed, past studies argue that low-latitude wind forcing associated with the Madden-Julian oscillation (MJO) and phases of the monsoon excite wave responses that effect intraseasonal sea-level variability along Sumatra and Java (Iskandar et al., 2005), the Bay of Bengal (Cheng et al., 2013), and India and Sri Lanka (Suresh et al., 2013; Dhage and Strub, 2016). Our results reinforce these past findings, and suggest that these nonlocal forcing effects mediated by large-scale wave responses continue on and are communicated to the Persian Gulf.

We perform a similar analysis with GRACE data. Correlations between ζ_0 and either GRACE R_m or its Hilbert transform $\mathcal{H}(R_m)$ over the Indian Ocean are shown in Figure 8. While there is essentially no meaningful correlation anywhere between ζ_0 and $\mathcal{H}(R_m)$, there is significant correlation between ζ_0 and GRACE R_m broadly over much of the Indian Ocean (Figure 8). This suggests that ζ_0 is also related to a basin-scale equilibrium response in addition to the more transient wave adjustments trapped to the coast and the equator suggested by the altimetry data (Figure 7). Indeed, the correlation pattern between ζ_0 and R_m (Figure 8a) is similar to the spatial structure of the intraseasonal fluctuation of the Indian Ocean identified by Rohith et al. (2019) based on data from bottom-pressure recorders, GRACE, and a general circulation model. They argue that wind-curl fluctuations at 30–80-day periods over the Wharton basin associated with the MJO excite planetary and topographic Rossby wave responses that lead to a basin-wide barotropic variation that is confined to the Indian Ocean by bathymetric contours. Our results provide observational evidence that this large-scale intraseasonal fluctuation affects variability not only over the deep Indian Ocean but also within its shallow marginal seas.

Wave propagation apparent in Figure 7 hints that ζ_0 variability may be predictable to some extent. That is, armed with upstream ζ information, it may be possible to anticipate ζ_0 variance in advance.

421 To test this possibility, we compute lagged correlation coefficients between ζ_0 and ζ at earlier times
 422 over the Indian Ocean. Results are shown in Figure 9 for lead times of 1 and 2 months. Considering
 423 a 1-month lead time, we find positive correlations between ζ_0 and ζ upstream along the Indian
 424 Subcontinent and Maritime Continent, from eastern India to Sumatra, and negative correlations
 425 over the western Equatorial Indian Ocean between Somalia and the Maldives (figure 9a). Indeed,
 426 the pattern of correlation between ζ_0 and ζ 1 month earlier is similar to the structure of correlation
 427 between ζ_0 and $\mathcal{H}(\zeta)$ (cf. Figures 7b, 9a), suggesting a dominant timescale of ~ 4 months. Values
 428 of 0.4–0.5 are apparent off Myanmar and Sumatra (Figure 9a), hinting that 16–25% of the variance
 429 in ζ_0 can be predicted from ζ knowledge in these regions 1 month earlier. Considering a lead time
 430 of 2 months, we observe that ζ_0 and ζ are largely uncorrelated, except for along Pakistan, western
 431 India, and Sri Lanka, where negative coefficients between -0.3 and -0.4 are seen. This implies that
 432 9–16% of the ζ_0 variance can be predicted from ζ observations along this coastline 2 months earlier.
 433 Considering lead times of 3 months and longer, we detect no significant correlations between ζ_0
 434 and ζ elsewhere (not shown), indicating that there is little skill in predictions of intraseasonal ζ_0
 435 variability more than 2 months into the future from wave characteristics and ocean memory alone.

436 **5. Summary and discussion**

437 We studied intraseasonal variability in ocean dynamic sea level (ζ) over the Persian Gulf during
 438 2002–2015 using satellite observations and other data (Figures 1, 2). Intraseasonal ζ variability
 439 in the Persian Gulf manifests in a basin-wide, vertically coherent mode of fluctuation (Figures 3–
 440 5). This large-scale mode is related to freshwater flux and barometric pressure over the Persian
 441 Gulf, wind stress along the Strait of Hormuz, and nonlocal forcing embodied in ζ variations at
 442 the boundary in the Gulf of Oman (Figure 6). The ζ boundary condition shows rich correlation
 443 patterns with altimetry data upstream along the Indian Subcontinent, Maritime Continent, and

equatorial Indian Ocean (Figure 7), and with GRACE data broadly over the Indian Ocean (Figure 8), suggesting an intimate connection between intraseasonal ζ variability in the Persian Gulf and large-scale circulation and climate over the Indian Ocean mediated by equatorial-, Rossby-, and coastal-wave processes identified previously (Cheng et al., 2013; Dhage and Strub, 2016; Iskandar et al., 2005; Oliver and Thompson, 2010; Rohith et al., 2019; Suresh et al., 2013, 2016; Waliser et al., 2003, 2004). Our results indicate that some intraseasonal ζ variance in the Persian Gulf may be predictable a month or so in advance from upstream observations and the physics of coastal wave propagation and ocean memory (Figure 9).

Our results establish the dominant magnitudes, scales, and mechanisms of intraseasonal sea-level variability in the Persian Gulf, and thus build on findings from past works that emphasize seasonal cycles and decadal trends (Al-Subhi, 2010; Alothman et al., 2014; Ayhan, 2020; El-Gindy, 1991; El-Gindy and Eid, 1997; Hassanzadeh et al., 2007; Hosseinibalam et al., 2007; Sharaf El Din, 1990; Siddig et al., 2019; Sultan et al., 1995a, 2000). Our study demonstrates that GRACE satellite retrievals are informative for interrogating coastal sea level over a semi-enclosed marginal sea, thereby complementing previous efforts that demonstrate the value of GRACE data in other marginal seas (Feng et al., 2012, 2014; Fenoglio-Marc et al., 2006, 2012; Landerer and Volkov, 2013; Loomis and Luthcke, 2017; Piecuch and Ponte, 2015; Piecuch et al., 2018; Tregoning et al., 2008; Wahr et al., 2014; Wang et al., 2015; Wouters and Chambers, 2010), and encouraging further exploration of GRACE data in the Persian Gulf at other timescales.

This investigation advances knowledge of sea-level variability in the Persian Gulf. It also paves the way for future studies, pointing to open questions. For example, we developed and tested a theory for a horizontally uniform fluctuation of the Persian Gulf. However, the leading mode of intraseasonal ζ variability in the region exhibits spatial structure, such that magnitudes are larger in the northwest and smaller in the southeast of the Persian Gulf (Figures 3, 5). We hypothesized that

468 this spatial structure could arise from local forcing by river runoff or wind stress over the Persian
469 Gulf. Future studies based on high-resolution ocean models should test this hypothesis and identify
470 the controls on spatial structure. It also remains to quantify whether baroclinic effects and steric
471 processes contribute to the dominant intraseasonal ζ variability in the Persian Gulf. Future studies
472 could explore this topic by comparing differences between altimetry and GRACE data, which are
473 potentially informative of steric processes, to sea-level changes expected from a passive response
474 to local surface heat flux (e.g., Cabanes et al., 2006).

475 We determined that dynamic response to barometric pressure and freshwater flux is a secondary
476 but nevertheless significant contributor to intraseasonal ζ variability in the Persian Gulf (Figure 6).
477 This is interesting, given that the barotropic ocean response to surface loading is generally expected
478 to be isostatic on timescales longer than a few days (e.g., Wunsch and Stammer, 1997; Ponte, 2006).
479 In our model physics, the dynamic response is permitted by friction through the Strait of Hormuz.
480 Our finding that freshwater flux elicits a ζ response on the order of a few mm (Figure 6) is consistent
481 with the basic ζ magnitudes simulated for this region across subdaily to annual timescales by Ponte
482 (2006) using a 1-year simulation from a global barotropic ocean general circulation model forced
483 with evaporation and precipitation (Hirose et al., 2001); however, that model was designed for
484 global studies, and it used coarse resolution ($\sim 1^\circ$) and a large friction coefficient ($2 \times 10^{-2} \text{ m s}^{-1}$),
485 which may not accurately capture important physics in and around the Persian Gulf. Future studies
486 using high-resolution ocean models would be informative for clarifying the nature of intraseasonal
487 ζ variation in the Persian Gulf and the role of surface loading. Also relevant here is the fact that the
488 non-isostatic response to barometric pressure is roughly in quadrature with the forcing (Table 4).
489 This highlights the importance of considering phase information when testing for departures from
490 a pure inverted-barometer response in sea-level data (e.g., Mathers and Woodworth, 2001, 2004).

Past studies argue that low-latitude wind forcing of the Indian Ocean related to large-scale climate modes excites wave responses that effect intraseasonal sea-level variability along the Indian Subcontinent and Maritime Continent, from Sumatra to western India (Cheng et al., 2013; Dhage and Strub, 2016; Iskandar et al., 2005; Suresh et al., 2013). We provide evidence that these coastal-trapped waves continue propagating downstream and influence sea level in the Gulf of Oman and Persian Gulf (Figure 7). We acknowledge that, while it suggests wave propagation, Figure 7 could alternatively indicate the spatial scales of the atmospheric forcing. For example, large-scale wind forcing along the equator and off the southern tip of the Indian subcontinent could simultaneously excite equatorial waves and coastal waves propagating in the cyclonic sense along the west coast of the Indian subcontinent (e.g., Suresh et al., 2013; Dhage and Strub, 2016). Future studies should identify the dominant centers of action of atmospheric forcing of intraseasonal ζ variability in the Persian Gulf, and whether coastal-trapped waves arriving in the Gulf of Oman have their origin in equatorial waves that impinged on the Maritime Continent. Our results also raise questions of whether such wave signals are felt even farther downstream along the coastal waveguide, for example, in the Red Sea. Previous investigations of sea-level variability in the Red Sea on timescales from days to decades largely emphasize the role of more local forcing (Abdelrahman, 1997; Churchill et al., 2018; Cromwell and Smeed, 1998; Osman, 1984; Patzert, 1974; Sofianos and Johns, 2001; Sultan and Elghribi 2003; Sultan et al., 1995b, 1995c, 1996). However, recent work by Alawad et al. (2017, 2019) suggests that mean sea-level variability in the Red Sea is partly related to large-scale modes of climate variability. These authors reason that this relationship is mediated by westward propagation of off-equatorial Rossby waves originating in the eastern tropical Indian Ocean. Based on our results, we hypothesize that coastal-wave propagation may also play a role in facilitating this relationship between sea level in the Red Sea and large-scale climate. We leave it to future studies to test this hypothesis.

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517 *Data availability statement.* Data are available through links provided in Table 1. Matlab codes
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519 upon request.

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Data set	Location
Altimetry	ftp://anon-ftp.ceda.ac.uk/neodc/esacci/sea_level/data/L4/MSLA/v2.0/
GRACE	https://podaac.jpl.nasa.gov/dataset/TELLUS_GRAC-GRFO_MASCON_CRI_GRID_RL06_V2
Tide gauges	https://www.psmsl.org/data/obtaining/complete.php
ERA-Interim	http://cmip5.who.edu/?page_id=566
GPCP	https://psl.noaa.gov/data/gridded/data.gpcp.html
OAFlux	ftp://ftp.who.edu/pub/science/oaflux/data_v3/monthly/evaporation/
JRA55-do	http://amaterasu.ees.hokudai.ac.jp/~tsujino/JRA55-do-suppl/runoff/

TABLE 1. Data sources. All websites are current as of this writing.

Station Name	Nation	PSMSL Identifier	Longitude (°E)	Latitude (°N)	Span	Completeness
Mina Sulman	Bahrain	1494	50.6	26.2	1979–2006	66.1%
Emam Hassan*	Iran	1868	50.3	29.8	1995–2006	91.7%
Bushehr*	Iran	1939	50.8	28.9	2004–2006	100.0%
Kangan*	Iran	1869	52.1	27.8	1995–2006	98.6%
Shahid Rajaei*	Iran	1870	56.1	27.1	1995–2006	100.0%

TABLE 2. Description of tide-gauge records. Asterisk indicates metric data without complete datum histories.

Parameter	Description	Value
ζ	Ocean Dynamic Sea Level	—
τ	Mean Wind Stress Along Strait of Hormuz	—
q	Surface Freshwater Flux	—
p	Barometric Pressure	—
ζ_0	Ocean Dynamic Sea Level in Gulf of Oman	—
$\bar{\tau}$	Spatial Average over Persian Gulf	—
S	Surface Area of Persian Gulf	$2.2 \times 10^5 \text{ km}^2$
H	Average Depth of Persian Gulf	30 m
L	Length of Strait of Hormuz	500 km
W	Width of Strait of Hormuz	100 km
g	Gravitational Acceleration	9.81 m s^{-2}
ρ	Ocean Density	1029 kg m^{-3}
r	Friction Coefficient [†]	$1 \times 10^{-3} - 1 \times 10^{-2} \text{ m s}^{-1}$
σ	Inverse Resonance Timescale	$1.8 \times 10^{-5} \text{ s}^{-1}$
λ	Inverse Frictional Timescale	$3.3 \times 10^{-5} - 3.3 \times 10^{-4} \text{ s}^{-1}$

TABLE 3. Descriptions of and, where applicable, reasonable values for variables and parameters in governing equations. [†]Values of the friction coefficient r are uncertain. Previous studies variously use values ranging from as small as $4 \times 10^{-5} \text{ m s}^{-1}$ (e.g., Ponte, 1994) to as large as $2 \times 10^{-2} \text{ m s}^{-1}$ (e.g., Ponte, 2006). Values in the table represent a reasonable, physically plausible range based on choices made in previous studies.

Parameter (Units)	Theoretical Range	Empirical Value
z_{ζ_0} (unitless)	0.8–1.0	1.0 ± 0.2
θ_{ζ_0} (degrees)	5–38	5 ± 10
z_{τ} (m Pa ⁻¹)	1.0–1.3	1.5 ± 0.5
θ_{τ} (degrees)	5–38	30 ± 25
$z_{\bar{q}}$ (days)	1.2–9.0	9.4 ± 3.7
$\theta_{\bar{q}}$ (degrees)	3–38	30 ± 27
$z_{\bar{p}}$ (cm mb ⁻¹)	0.1–0.5	0.8 ± 0.5
$\theta_{\bar{p}}$ (degrees)	56–87	65 ± 52

857 TABLE 4. Estimates of the scaling coefficients (z_j) and phase angles (θ_j) in Eq. (4). The theoretical ranges are
 858 determined by averaging Eqs. (5)–(12) over the range $\omega = 2\pi / (6 \text{ months})$ to $2\pi / (2 \text{ months})$ using the constant
 859 values for σ , L , ρ , g , and H and the minimum and maximum values for λ tabulated in Table 3. Empirical values
 860 are determined through multiple linear regression involving $\bar{\zeta}$ and ζ_0 from altimetry, τ and \bar{p} from ERA-Interim,
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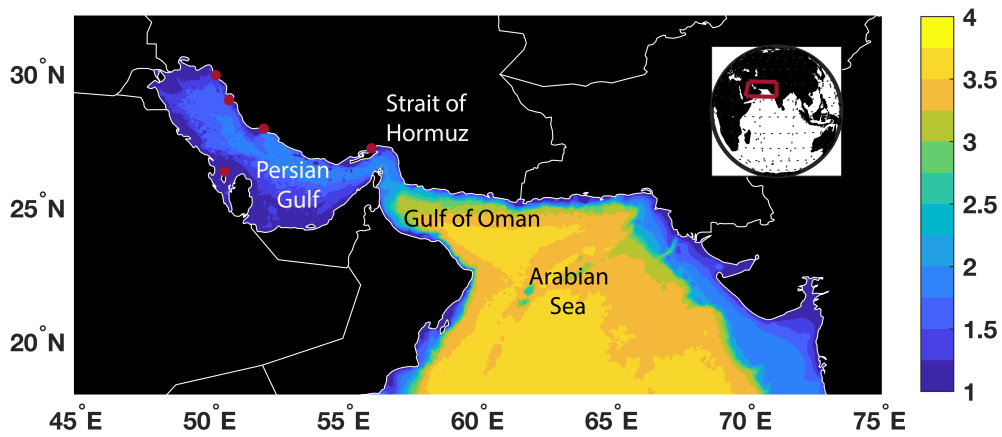


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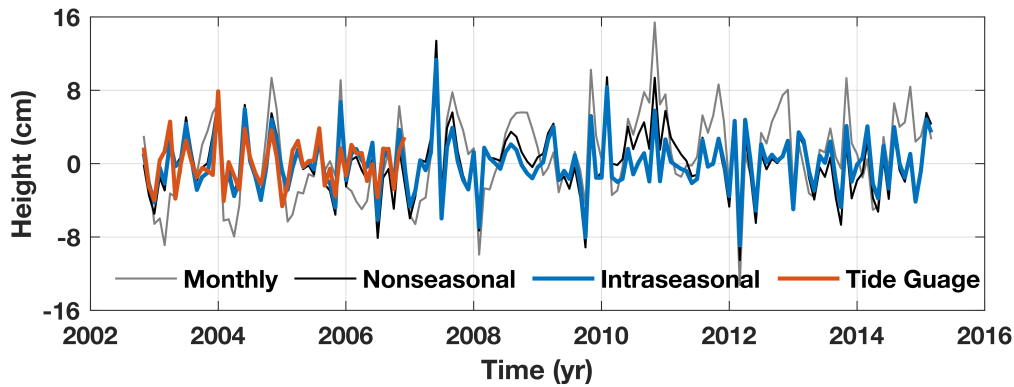


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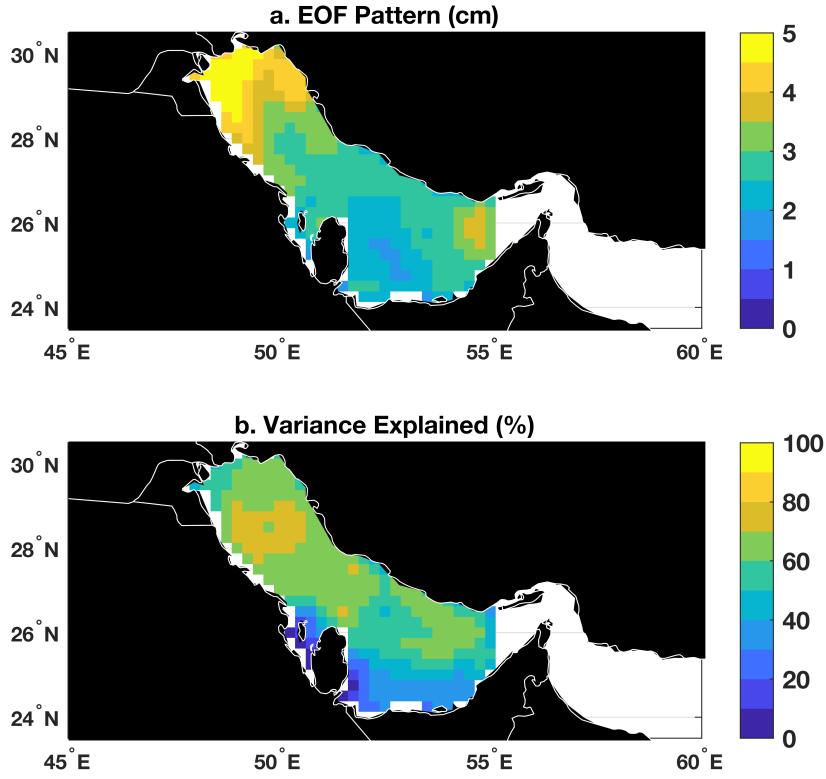


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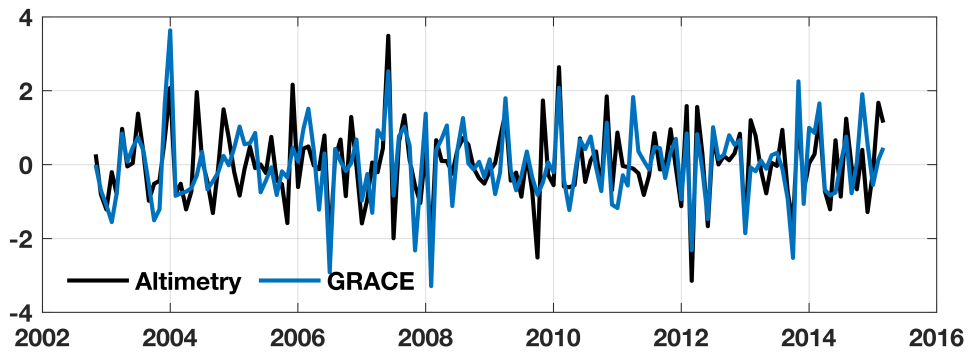


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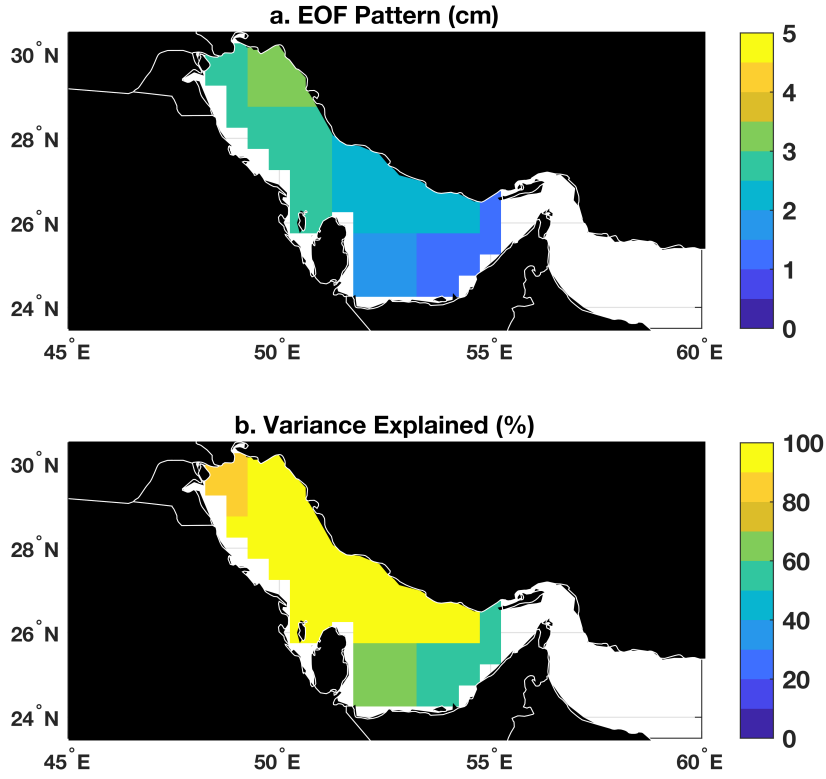


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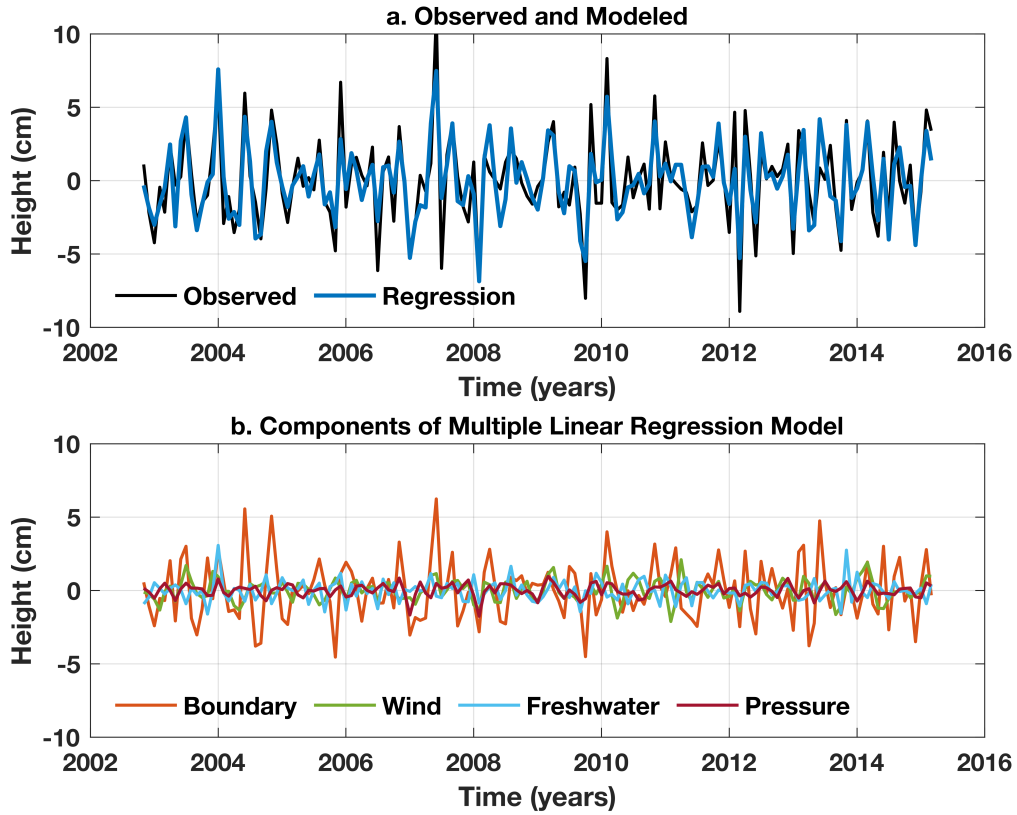


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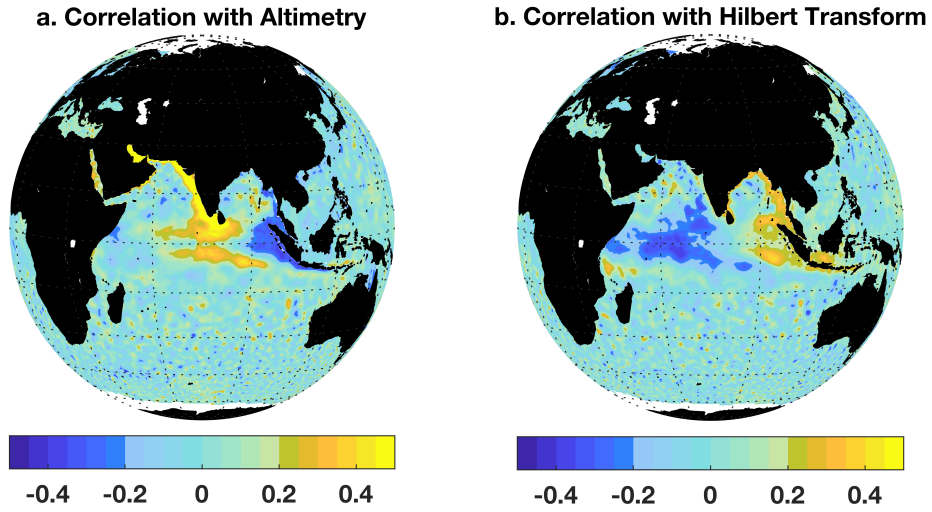


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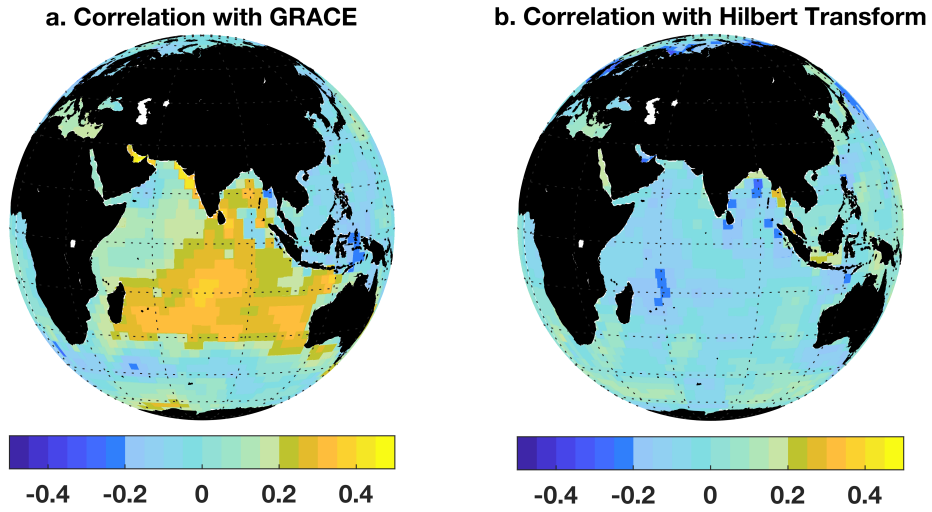


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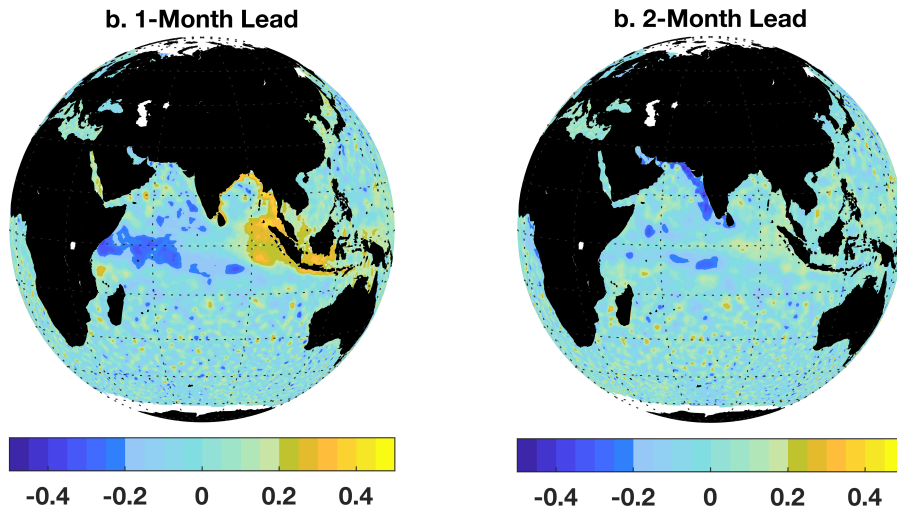


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