

1           **ANCHOR: Global Parametrized Ionospheric Data  
2           Assimilation**

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9           **Distribution Statement**

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11          **Key Points:**

- 12          • Novel approach for rapid ionospheric data assimilation using anchor points  
13          • Global ionospheric data assimilation for plasma density parameters  
14          • Extraction of the anchor points from the ionospheric data

15      **Abstract**

16      ANCHOR is a novel assimilative model developed at the U.S. Naval Research Laboratory.  
 17      It extracts ionospheric parameters from RO and ionosonde data and assimilates  
 18      them as point measurements into the maps of the background parameters using Kalman  
 19      Filter approach. This paper introduces the ANCHOR algorithm, discusses its coordi-  
 20      nate system and background, explains the background covariance formation, discusses  
 21      the extraction of the ionospheric parameters from the data and the assimilation process,  
 22      and, finally, shows the results of the observing system simulation experiment.

23      **1 Introduction**

24      The ionosphere starts  $\sim$ 80 km from the Earth's surface and extends all of the way  
 25      to the exosphere, which is thousands of kilometers above the Earth's surface. The free  
 26      electrons that inhabit the ionosphere refract the electromagnetic waves that traverse through  
 27      the ionosphere. Since high frequency (HF) communication is made possible by the prop-  
 28      agation of electromagnetic waves through the ionosphere, it is crucial to know the amount  
 29      of electrons along the communication path to establish the HF communication link by  
 30      choosing the best transmission frequency. Climatological models, such as the Interna-  
 31      tional Reference Ionosphere (IRI) (Bilitza et al., 2022; Forsythe et al., 2023), provide good  
 32      estimates of the electron density for a given time and solar activity level. However, the  
 33      ionospheric state often deviates from the climatological prediction because it is a very  
 34      active environment with localized structures and short-lived events. One way around this  
 35      problem is to assimilate the ionospheric observations into the model background in or-  
 36      der to obtain a realistic specification.

37      There are many ionospheric data assimilation (DA) models that use a Kalman fil-  
 38      tering (KF) approach, including Global Assimilation of Ionospheric Measurements (GAIM)  
 39      (Schunk et al., 2004), the Electron Density Assimilative Model (EDAM) (Angling & Can-  
 40      non, 2004), the Ionospheric Data Assimilation Four-Dimensional (IDA4D) model (Bust  
 41      & Crowley, 2007), and a Three-Dimensional Regional Assimilative Model of the Iono-  
 42      spheric Electron Density (Pignalberi, 2019). As a separate example, the IRI-Based Real-  
 43      Time Assimilative Mapping (IRTAM) (Galkin et al., 2020) DA model uses KF on the  
 44      time series of the ionospheric parameters, and is further discussed in Section 4.2.

45 There are several known disadvantages related to the non-parametrized DA mod-  
46 els. First, a high computational load is required when performing the calculations for  
47 the 3-D electron density grid. When a global regularly-spaced grid with 2° horizontal res-  
48 olution (or 16200 horizontal grid points) and 10 km vertical resolution (or 100 vertical  
49 grid points) is used, the total number of grid points becomes 1,620,000. The background  
50 covariance matrix for this grid should then be as large as  $1,620,000 \times 1,620,000$ . This  
51 matrix size slows down the calculations significantly.

52 The second challenge relates to the ingestion of non-local data types, such as slant  
53 total electron content (sTEC) from radio occultation (RO) data, sTEC from ground-based  
54 Global Positioning System (GPS) receivers, and ultraviolet (UV) radiance satellite imag-  
55 ing data. In the case of RO data, the sTEC measurement indicates the amount of elec-  
56 trons along a straight line between Low-Earth Orbit (LEO) and GPS satellites. How-  
57 ever, the DA system has no information on how the data should be distributed along this  
58 path, or at which points along this straight line the modeled ionosphere should be cor-  
59 rected to match the observations. This makes the sTEC data non-localizable, with a very  
60 low information content on a per-measurement basis (Reid et al., 2023). One consequence  
61 of the ingestion of non-local data is difficulty resolving the horizontal variation of the elec-  
62 tron density peak height (or  $hmF2$ ) parameter. For example, the ingestion of RO sTEC  
63 into IDA4D mainly influences the peak density ( $NmF2$ ), and infrequently adjusts the  
64  $hmF2$  (Forsythe & McDonald, 2022).

65 The third challenge relates to the ingestion of sTEC measurements in data-dense  
66 regions. As was previously noted by Forsythe et al. (2021), the intersection of sTEC rays  
67 can cause localized changes to electron density profiles (EDPs), creating unrealistic EDPs  
68 with fictitious layers or bite-outs. In other words, when sTEC rays intersect in an un-  
69 determined way, it is possible for the assimilation to produce the analysis that matches  
70 the measurements well, but contains non-physical EDPs. This last challenge is of par-  
71 ticular importance for the DA applications, such as ray tracing, because an unrealistic  
72 bottom-side electron density will result in incorrect HF signal propagation path that lead  
73 to unusable communication recommendations.

74 These challenges provided the motivation to rethink the current approach to iono-  
75 spheric DA. This led to the development of ANCHOR, a novel DA approach that assim-  
76 ilates anchor points extracted from data into horizontal maps of the ionospheric param-

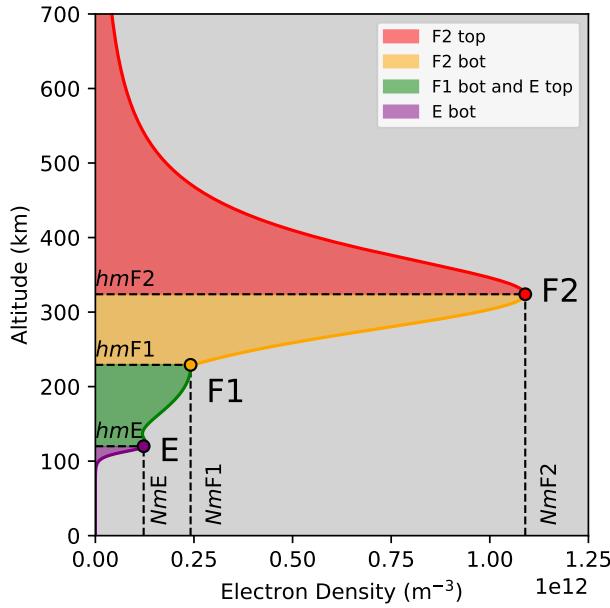
eters. Naval Research Laboratory (NRL) currently implements this DA method in Python, as it has low computational requirements.

This paper starts by discussing the parametrization method used in ANCHOR. It then describes the coordinate system, background model (PyIRI), derivation of the background covariance, and DA scheme used in ANCHOR. Finally, it shows the results of an observing system simulation experiment (OSSE) that evaluates ANCHOR's performance.

## 2 ANCHOR DA

### 2.1 Parametrization

According to the IRI model (Bilitza et al., 2022), a typical vertical ionospheric profile consists of a main layer, called the F2 layer, that peaks at approximately 350 km. Figure 1 demonstrates a typical EDP. During the daytime, the photochemical layers (F1 and E) appear below the F2 peak. The D layer, which overlaps with the upper mesosphere is neglected in PyIRI due to its low plasma density. The D region is also ignored by ANCHOR due to the lack of D-region data to assimilate.



**Figure 1.** Example of an EDP and its main parameters.

It is possible to describe the EDP in terms of 11 parameters, where 6 parameters  $NmF2$ ,  $NmF1$ ,  $NmE$ ,  $hmF2$ ,  $hmF1$ , and  $hmE$  indicate the peaks and heights of the three layers, that are shown with red, orange and purple circles, respectively, in Figure 1. The remaining, five parameters are needed to describe the thicknesses of the top and the bottom sides of F2 layer ( $B_{top}^{F2}$ ,  $B_{bot}^{F2}$ ), the bottom side of F1 layer ( $B_{bot}^{F1}$ ) and the bottom side of E layer ( $B_{bot}^E$ ), as shown with colors in Figure 1. The 11 parameters can be reduced to five independent anchor points (hence the name ANCHOR):  $NmF2$ ,  $hmF2$ ,  $B_{top}^{F2}$ ,  $B_{bot}^{F2}$ , and  $NmE$ , as the rest of the parameters can be either fixed at a constant values, such as  $hmE$ ,  $B_{bot}^E$ , and  $B_{top}^E$ , or can be derived from the combination of the independent parameters, as it is done in PyIRI for the  $NmF1$ ,  $hmF1$ , and  $B_{bot}^{F1}$  (Forsythe et al., 2023). Although, ANCHOR currently uses five independent anchor parameters, it is possible to extend this number in the future (e.g. to include the  $hmE$  correction, the D-region parameters, and the sporadic E layer).

## 2.2 Reference Frame and The Background

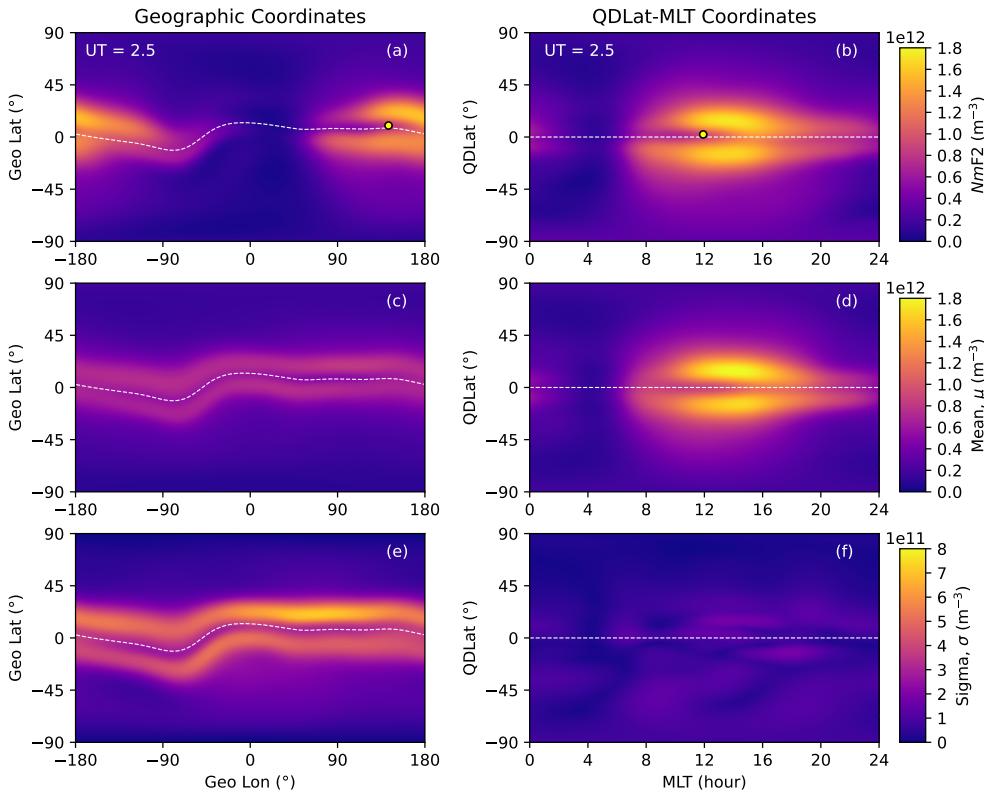
A reference frame selection is a very important step in the DA set-up process. Most ionospheric DA models operate in geographic coordinates. In this frame the subsolar point changes longitude position as the Earth rotates, whereas the position of the geomagnetic equator stays unchanged with time. As a result, the position of the ionospheric regions (in relation to the Earth's surface) also change with time. Figure 2a shows a snapshot of the  $NmF2$  according to PyIRI for 02:30 Universal Time (UT) of 1 April 2020. The yellow circle shows the position of the subsolar point. The two crests of the Equatorial Ionization Anomaly (EIA) flank the geomagnetic equator shown with white dashed line.

A collection of the 24-hour snapshots in geographic coordinates describes the complete cycle of “motion” of the ionosphere from the right to the left (as the Earth rotates). Therefore, the daily mean  $\mu$  and the standard deviation  $\sigma$ , shown in Figures 2c and 2e, respectively, among these snapshots determined in the geographic coordinates look as equally distributed in longitude enhancements that follow the curvature of the geomagnetic equator.

However,  $\mu$  and  $\sigma$  parameters look completely different if calculated in another coordinate system. Figure 2b shows the  $NmF2$  snapshot (same as in Figure 2a) organized in the Quasi-Dipole Latitude (QDLat) and Magnetic Local Time (MLT) coordinate sys-

123 tem. The subsolat point is now centered at 12 MLT and the geomagnetic equator is a  
 124 straight line. The two crests of the EIA flank the straight geomagnetic equator.

125 Importantly, the other UT snapshots in QDLat-MLT coordinates look very sim-  
 126 ilar to Figure 2b. Therefore, the daily mean  $\mu$ , shown in Figure 2d, also looks very sim-  
 127 ilar to each UT snapshot. Since the variance between each UT snapshot is very small,  
 128 the standard deviation ( $\sigma$ ) in QDLat-MLT coordinates is very low, especially in com-  
 129 parison to  $\sigma$  from Figure 2e.



**Figure 2.** Comparison of the  $NmF2$  and its daily variability expressed in geographic and QDLat-MLT coordinate systems for 1 April 2020. Panels (a) and (b) show the snapshots at 02:30 UT. Panels (c) and (d) show the daily mean, and panels (e) and (f) show the daily standard deviation. Geomagnetic equator is shown with white dashed line. Yellow circles in panels (a) and (b) show the locations of the subsolar points.

130 ANCHOR performs DA in the QDLat-MLT coordinate system. Unlike other DA  
 131 models that use geographic coordinate systems that are agnostic to the subsolar loca-  
 132 tion, the QDLat-MLT coordinate system considers both the orientation of the magnetic

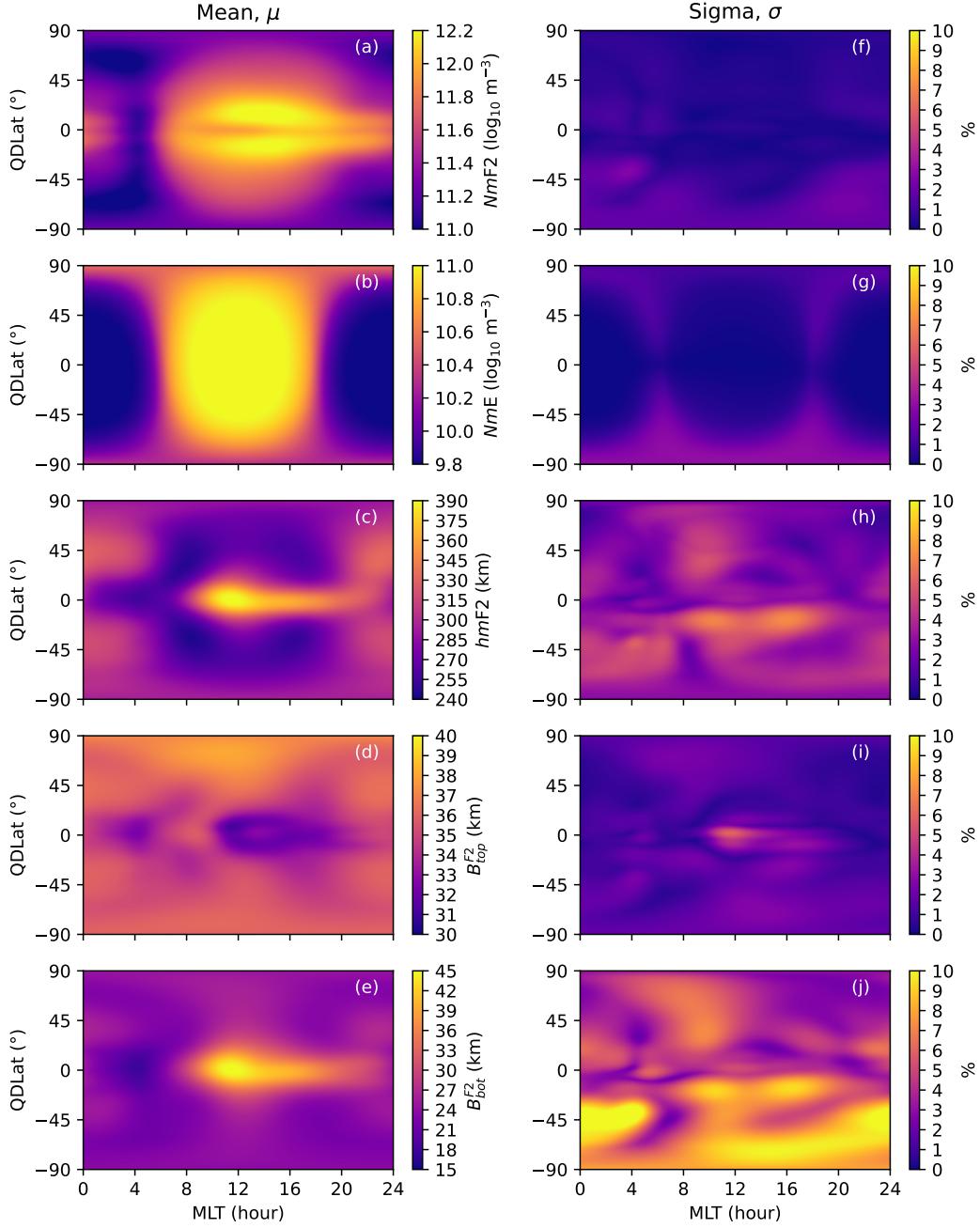
133 field and the subsolar point location. Since this information is essential in the organi-  
 134 zation of ionospheric density (through photoionization and plasma transport), this ref-  
 135 erence frame is a more natural fit for ionospheric specification. This is reflected in the  
 136 consistency of  $\mu$  over the course of a day in the QDLat-MLT frame. The model variance  
 137 and covariance also need to be defined specifically for QDLat-MLT system and are de-  
 138 scribed below.

139 At the current stage of development, the PyIRI model (Forsythe et al., 2023) is em-  
 140 ployed as a background for ANCHOR. The background ionospheric parameters are de-  
 141 termined for 15-min time frames because the regularly-spaced in QDLat-MLT coordi-  
 142 nates grid points have different geographic locations in a geographic coordinate system  
 143 (the PyIRI input needs to be in the geographic coordinate system). The  $\mu$  and  $\sigma$  for all  
 144 PyIRI background parameters for 1 April 2020 are shown in Figure 3.  $\sigma$  is expressed as  
 145 percent difference from the  $\mu$ . ANCHOR operates in  $\log_{10}$  space for the electron den-  
 146 sity peaks to guarantee the positivity of the outputs. Therefore,  $NmF2$  and  $NmE$  pa-  
 147 rameters in Figure 3a and 3b are expressed in  $\log_{10}$  scale. The  $NmE$  parameter is strongly  
 148 controlled by the solar irradiance, while the other parameters show spatial variations that  
 149 reveal the importance of plasma transport. All of the parameters have very low diurnal  
 150 standard deviations. When expressed in percentage of the mean, the density anchor pa-  
 151 rameters  $NmF2$  and  $NmE$  have standard deviations that do not exceed two percent, with  
 152 the largest standard deviations in the Antarctic region, where observations are sparse.  
 153 The percent standard deviation is most significant for the  $B_{bot}^{F2}$  parameter.

154 The model daily  $\sigma$ , as is, cannot be used as model covariance  $\sigma$ , for the following  
 155 reasons. First, it is so low that it would not allow the background to deviate far enough  
 156 to match the observations. Second, the horizontal structure of  $\sigma$  should be considered  
 157 with caution. For instance,  $\sigma$  for the  $NmE$  is equal to zero around noon and midnight,  
 158 as shown in Figure 3g, because the E-region model takes into account only the position  
 159 of the sun. These zero values for  $\sigma$  will prevent the propagation of the information from  
 160 the data to the model, even if the variance is inflated using some multiplier. Another ex-  
 161 ample can be made about the  $NmF2$   $\sigma$ , shown in Figure 3f. During the day time, it has  
 162 zero values in the Northern hemisphere. This means that the data observed in the North-  
 163 ern hemisphere will have minimum DA influence compared to the data observed in the  
 164 Southern hemisphere, just because the background will be restricted to move away from  
 165 the  $\mu$ .

166        The horizontal distribution of  $\mu$  is more important in the QDLat-MLT coordinate  
167        system, than the horizontal distribution of  $\sigma$ . For each parameter it clearly indicates where  
168        the values should be high, and where they should be low. Therefore, ANCHOR mod-  
169        els the variance, setting it to 30% of the daily background mean. It is important to men-  
170        tion that 30% is not important by itself as a number, because for the DA the ratio be-  
171        tween the variance of the model to the variance of the data is more important. For our  
172        OSSE (Section 3.1), the data error will be set to 10%, meaning that the DA should trust  
173        the data three times more than it trusts the background.

174        The background vector  $\vec{x}_b$  is formed for each time frame as a flattened 1-D array  
175        that contains all the background parameters. For example, in case of 15 min MLT and  
176        2° QDLat resolutions, the grid for each map of the parameters has the size of 8827 el-  
177        ements, and the  $\vec{x}_b$  vector has 44135 elements. A regular DA scheme (with similar hor-  
178        izontal resolution and 10 km vertical resolution) would have a state vector 183 times as  
179        large. This reduction in size allows much faster computation and lower hardware require-  
180        ments.



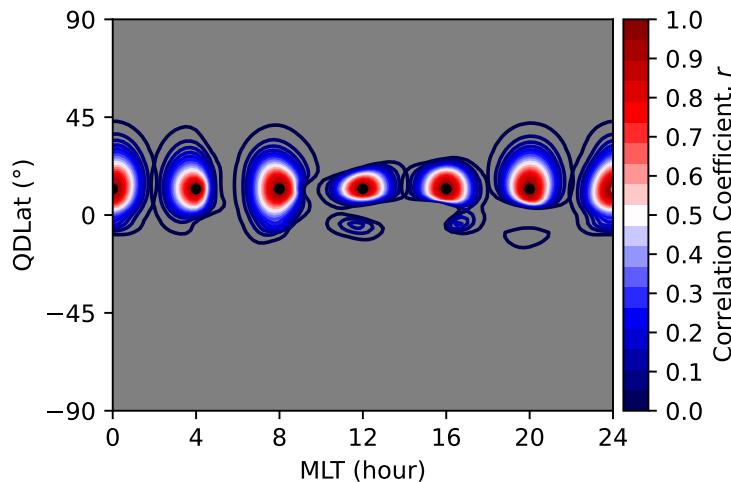
**Figure 3.** Daily mean  $\mu$  and standard deviation  $\sigma$  of the anchor parameters in QDLat-MLT coordinates for 1 April 2020.  $\sigma$  is expressed as percent from the  $\mu$ .

### 2.3 Background Covariance Matrix

Now that the model variance is established, it is necessary to develop the model covariance. A novel element here is the coordinate frame. ANCHOR calculates the back-

ground covariance,  $\tilde{P}_b$ , for each day of assimilation. First, it finds the deviation of each time frame from the daily mean for each anchor parameter. Then it finds the correlations of these deviations for each grid location with other locations. Since the correlations are calculated and not modeled as Gaussian distributions, it is necessary to localize them. Without localization one observation would transport the information all the way to the opposite side of the globe. Therefore, the correlations are multiplied with Gaspari-Cohn localization (Gaspari & Cohn, 1999) with the radius of  $20^\circ$  of great circle distance (GCD). This localization guarantees that at  $40^\circ$  GCD the correlation fully drops to zero. The localization radius of  $20^\circ$  GCD was chosen based on the previous work by Forsythe, Azeem, and Crowley (2020), where according to the GPS data the maximum of the mean horizontal ionospheric correlation length was found to be approximate  $20^\circ$  GCD.

The localized correlations for the  $NmF2$  parameter around six reference points located at  $12^\circ$  QDLat and separated by 4 MLT are shown in Figure 4. It is obvious that in QDLat-MLT coordinates the mean MLT position of the equatorial crests determines the unique shape of the correlation ellipses. Whereas three reference points at 12, 16, and 20 MLT have secondary correlation peaks that indicate the geomagnetically conjugate points on the opposite side of the equator that are present only for the  $NmF2$  parameter. This happens due to the latitudinal symmetry in the diurnal variance of the EIA. Importantly, the applied localization radius is large enough and does not suppress the horizontal structure of the correlations.



**Figure 4.** Correlations for  $NmF2$  around six reference points for 1 April 2020.

204        The background covariance matrix is formed for all of the anchor parameters, tak-  
 205        ing into account the localized correlations for the off-diagonal elements and the allowed  
 206        30% standard deviation for the diagonal elements. It is important to note that because  
 207        each parameter varies within this limit, it is impossible to obtain an unrealistic EDP. For  
 208        example, it is impossible to obtain a fictitious ionospheric layer or a bite out, because  
 209        it would require the addition of more parameters to the EDP description formalism. This  
 210        fully addresses the third challenge of the non-parametrized ionospheric DA listed in the  
 211        Section 1.

212        The covariance matrix is, however, further reduced to contain only the covariance  
 213        between the same parameters, excluding the cross-parameter correlations. Therefore, at  
 214        the current state of development, ANCHOR treats all five core parameters as indepen-  
 215        dent of each other. The cross-parameter correlation requires further investigation and  
 216        is beyond the scope of this study.

## 217        2.4 Observation Vector

218        Ionospheric parameters are extracted from RO and ionosonde data. This process  
 219        is described in the upcoming Section 3.1. The observation vector,  $\vec{y}$ , is formed as a 1-  
 220        D array of all the measured parameters, combining RO and ionosonde measurements,  
 221        following the same order of the parameters as in the  $\vec{x}_b$  array. The observation error co-  
 222        variance matrix,  $\tilde{R}$ , is formed assuming that the parameters are independent and set-  
 223        ting the observation error to 10% of the measurement. As was mentioned before, in com-  
 224        bination with the given 30% of the background variance, the DA trusts the data three  
 225        times more than it trusts the background.

## 226        2.5 The Observation Operator

227        The observation operator,  $\tilde{H}$ , is formed as a matrix with shape  $[N_{obs}, N_b]$ , where  
 228         $N_{obs}$  is the number of observations in array  $\vec{y}$ , and  $N_b$  is the number of elements in the  
 229         $\vec{x}_b$  array. The nearest-neighbor interpolation is used as a convenient simplification for  
 230        the localized point measurements. For each observation,  $\tilde{H}$  will have an element equal  
 231        to unity at the closest grid point to the occurred observation, and zeroes everywhere else.  
 232        Thus, the following matrix multiplication

$$\vec{y}_b = \tilde{H} \vec{x}_b \quad (1)$$

233 gives the expected observations  $\vec{y}_b$  according to the given background.

234 The  $\vec{y}_b$  is further used to exclude observations that differ significantly from the back-  
 235 ground. For the first three parameters ( $NmF2$ ,  $NmE$ , and  $hmF2$ ) the data points are  
 236 excluded in the case where the residuals  $|\vec{y} - \vec{y}_b|$  are larger than three times the back-  
 237 ground standard deviation. This is a so-called  $3\sigma$  rule. However, this multiplier is in-  
 238 creased to 10 for the  $B_{top}^{F2}$  and  $B_{bot}^{F2}$  since the residuals are significantly higher for those  
 239 parameters (as expected given the empirical model variation and/or worse performance  
 240 of the model for these parameters).

241 Additionally, it was found useful to combine observations that are close (located  
 242 within  $10^\circ$  GCD) and differ by more than 10%. The presence of such contradicting ob-  
 243 servations often creates an unrealistically sharp gradient in the analysis. Therefore, such  
 244 observations are averaged to reduce the number of contradicting measurements. The num-  
 245 ber of these exclusions is close to zero for the simulated data, but this reduction becomes  
 246 important for the real data (Sakov & Sandery, 2017).

## 247 2.6 Kalman Filter

248 The Kalman Filter formalism is used for the ANCHOR DA scheme, where the anal-  
 249 ysis  $\vec{x}_a$  represents the background corrected by the ingested observations and is defined  
 250 as

$$\vec{x}_a = \vec{x}_b + \tilde{K} [\vec{y} - \vec{y}_b], \quad (2)$$

251 where  $\tilde{K}$  is the Kalman gain defined as

$$\tilde{K} = \tilde{P}_b \tilde{H}^T \left[ \tilde{H} \tilde{P}_b \tilde{H}^T + \tilde{R} \right]^{-1}, \quad (3)$$

252 In such a configuration of the filter, the information from the previous time frame  
 253 does not propagate to the next time frame. For ANCHOR this is kept on purpose, be-  
 254 cause in the future ANCHOR will be coupled with full physics-based SAMI3 (Huba et  
 255 al., 2000) model. This model coupling will take care of the information propagation from  
 256 the current DA time frame to the next one.

257 However, for this paper ANCHOR will also be run in a different filter setting known  
 258 as Gauss-Markov model, where the  $\vec{x}_a$  is defined as

$$\vec{x}_a = \vec{x}_f + \tilde{K} [\vec{y} - \vec{y}_f], \quad (4)$$

259 where  $\vec{x}_f$  is the forecast calculated as

$$\vec{x}_f = \vec{x}_{a'} + (\vec{x}_{a'} - \vec{x}_b)e^{-dt/\tau}, \quad (5)$$

260 where  $\vec{x}_{a'}$  is the analysis from the previous assimilation time frame,  $dt$  is the DA time  
 261 resolution (0.25 hours), and  $\tau$  is decorrelation time that defines how soon the data be-  
 262 comes old (Forsythe, I. Azeem, et al., 2020), which is set to 2 hours. Appendix A has  
 263 more information on the decorrelation time for ANCHOR.

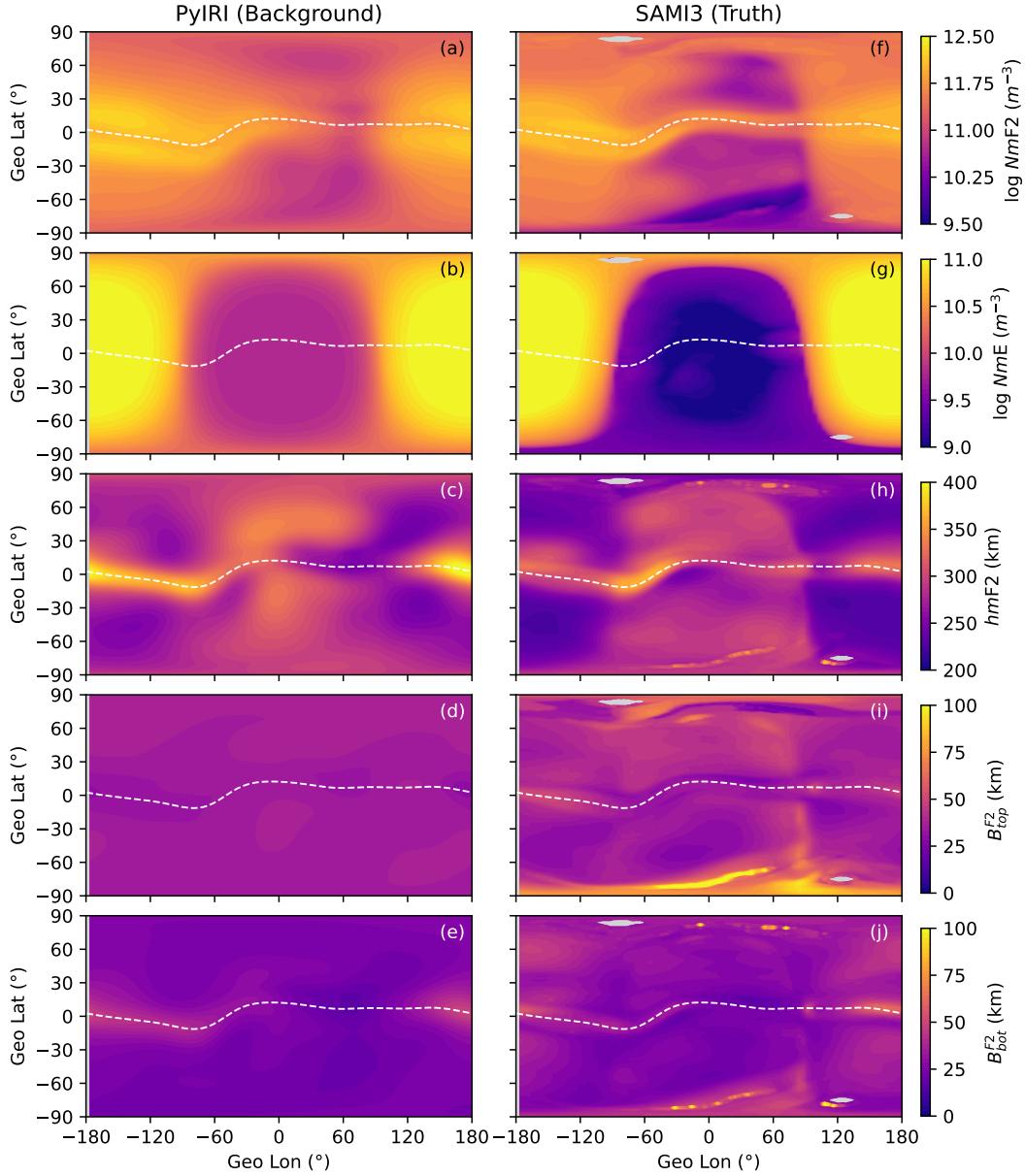
## 264 2.7 3-D Density

265 After each time frame ANCHOR finds other dependent parameters, such as  $NmF1$ ,  
 266  $hmF1$ , and  $B_{bot}^{F1}$ , utilizing the analysis of the main anchor parameters. Then, the maps  
 267 of all parameters are converted to 3-D electron density. ANCHOR uses the same formal-  
 268 ism as PyIRI to calculate the dependent parameters and to perform the computa-  
 269 tionally efficient 3-D density derivation, as described in details in Forsythe et al. (2023).

## 270 3 Results and Performance Evaluation

### 271 3.1 OSSE Truth and Simulated Observations

272 The evaluation of the ANCHOR algorithm was performed using an OSSE, set up  
 273 with the full physics-based SAMI3 model outputs as *truth*. Observations were simulated  
 274 based on the SAMI3 electron density output resampled into a regularly spaced geographic  
 275 grid. Figure 5 compares the PyIRI background parameters to the SAMI3-derived pa-  
 276 rameters. All the parameters besides the  $NmF2$  show a significant deviation from the  
 277 background, which is desirable for the DA evaluation.

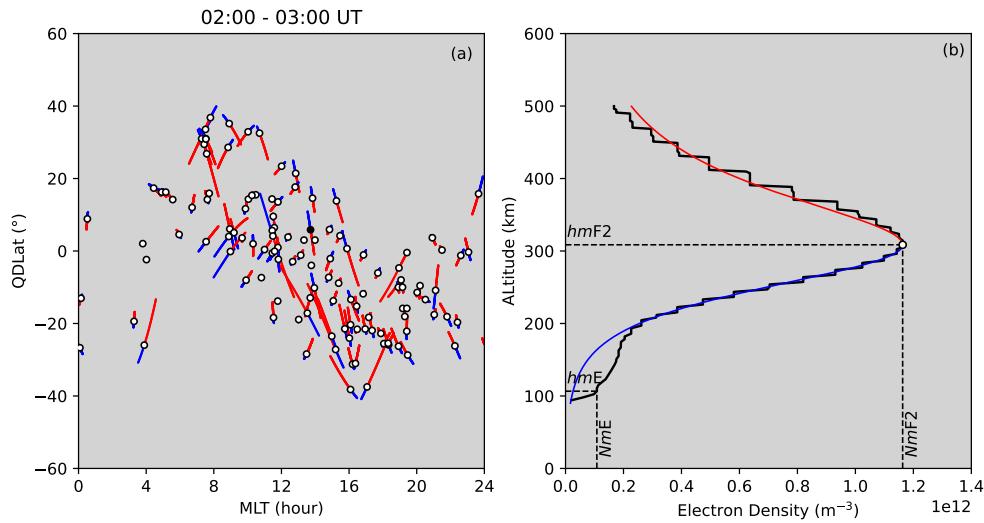


**Figure 5.** Comparison between background parameters and SAMI3 parameters used as known truth for 1 April 2020, 00:00:00 UT. Geomagnetic equator is shown with white dashed line.

COSMIC-2 (Lin et al., 2020) satellite EDPs along the tangent points (TP) were simulated using SAMI3 as a known truth for 1 April 2020 (by finding electron density along the TP locations from SAMI3). It is important to mention that this approach simplifies the OSSE, making the results look better than the real-life DA performance, where the RO sTEC measurements are first inverted to EDP using the Abel inversion technique. This choice is discussed in the Section 4. Figure 6a shows the locations of the COSMIC-

284 2 ROs between 2 and 3 UT. The white circles show the location of the observed  $NmF2$ ,  
 285 with red and blue lines indicating the locations of top- and bottom-sides of the F2 re-  
 286 gions along the TPs, respectively. A black circle in Figure 6a shows the EDP shown in  
 287 Figure 6b. The black line in Figure 6b shows the simulated profile, which is not smooth  
 288 because the vertical resolution of the TP is higher than the SAMI3 output file vertical  
 289 resolution. The  $NmF2$  and  $hmF2$  are first found. Then, using least square fitting, the  
 290 thicknesses parameters  $B_{top}^{F2}$  and  $B_{bot}^{F2}$  are found using Epstein function formalism that  
 291 is explained in details in Section 4 of the PyIRI paper (Forsythe et al., 2023). Red and  
 292 blue curves in Figure 6b demonstrate the reconstruction of the topside and bottomside  
 293 using the  $B_{top}^{F2}$  and  $B_{bot}^{F2}$  parameters, respectively. Next, the blue curve is subtracted from  
 294 the black curve and a location of the  $NmE$  is found as a local maxima around 110 km  
 295 with a 10 km allowed deviation.

296 The observation operator that was discussed in Section 2.5 localizes the RO mea-  
 297 surements to the location of the observed  $NmF2$ , or the white circles in Figure 6a. How-  
 298 ever, given the data influence radius defined by the background covariance model (shown  
 299 in Figure 4), the data influence will be distributed around the points, fully enclosing most  
 300 of the TP tracks in Figure 6a, especially the bottom side portions, shown with red color,  
 301 since they are usually shorter than the top-side portions.

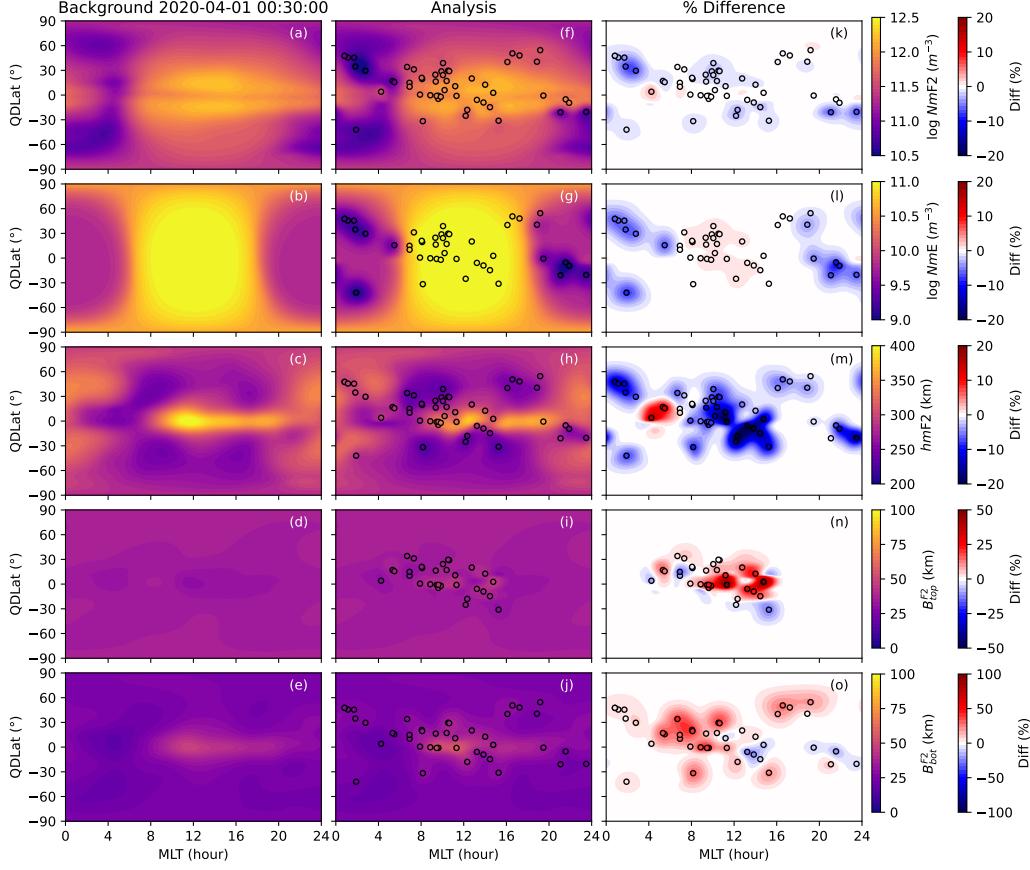


**Figure 6.** Locations of COSMIC-2 ROs and an example of the simulated data EDP.

302 A similar fitting procedure was applied to the vertical ionosonde data collected for  
303 1 April 2020 for 20 ionosondes. Additionally, the locations of the ionosondes were ran-  
304 domly perturbed to mask the operational data set. The parametrization differed from  
305 RO only in excluding the determination of the  $B_{top}^{F2}$ , since ionosondes can not measure  
306 the ionospheric density above the peak density. In this OSSE the simulated data rep-  
307 resent the ideal synthetic observation, without adding any additional instrumentation  
308 error.

309 **3.2 DA Analysis**

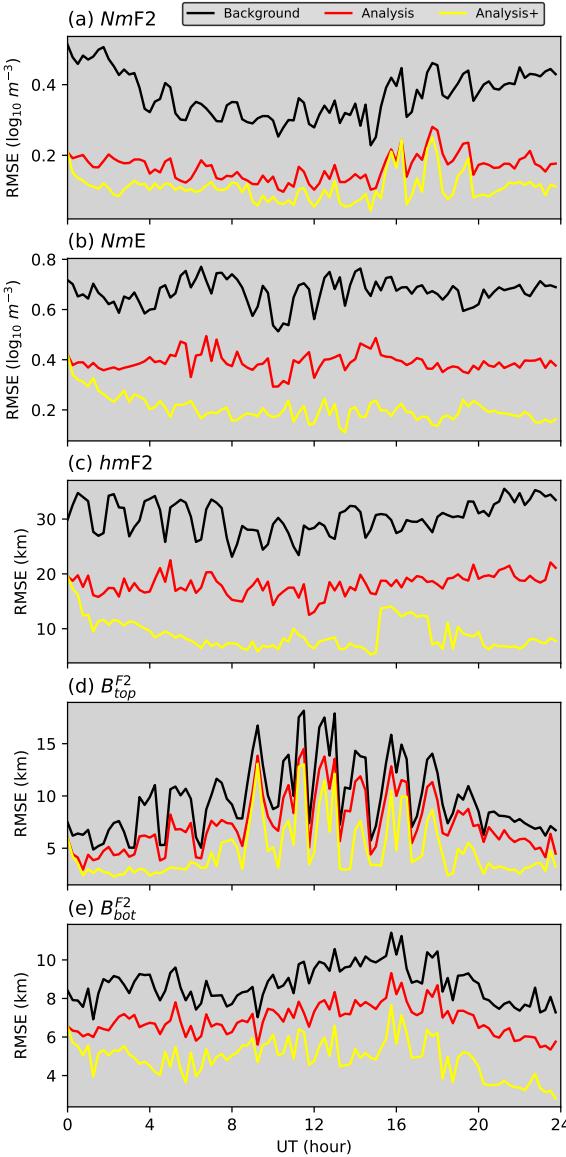
310 Figure 7 shows the results of the assimilation for one 15-min time frame for 1 April  
311 2020, 00:30:00 UT. The first column in Figure 7 shows the background parameters, the  
312 second column shows the analysis with the ingested point measurements shown in cir-  
313 cles, and the third column shows the percent difference between the background and the  
314 analysis together with the difference in the expected and observed parameters shown with  
315 circles. The largest modification to the background is made for the  $B_{top}^{F2}$  and  $B_{bot}^{F2}$  pa-  
316 rameters, which use expanded color bars to ensure all % differences are visible.



**Figure 7.** ANCHOR results for 1 April 2020, 00:30:00 UT.

### 3.3 Root Mean Square Errors

To analyze the ANCHOR results, the root mean square errors (RMSEs) between the SAMI3 truth and both the background and the analysis are found for the locations that have non-zero percent difference shown in the third column of Figure 7. Five panels in Figure 8 show the RMSEs for the background and analysis with black and red lines, respectively, for the five anchor parameters. In the case of  $NmF2$  and  $NmE$ , the RMSEs are reduced by half in the analysis. In case of  $hmF2$  the RMSEs are reduced by about 35%. Due to the large difference between the truth and the background, the  $B_{top}^{F2}$  and  $B_{bot}^{F2}$  show slightly lower reduction of RMSEs in the analysis.



**Figure 8.** RMSEs for two ANCHOR runs.

In case one wonders about the 1.5-hour oscillations in RMSEs, they arise due to  
 the clustering of the RO data. Figure 7i showed only RO data locations (since ionosonde  
 data does not have  $B_{top}^{F2}$  parameter). For this 15-min time window the RO points are clus-  
 tered between 4 and 16 MLT. Whereas the MLT sector moves by 4 MLT hours to the  
 right every 15 minutes, making the full loop in 1.5 hours. Since the RMSE calculation  
 includes only data-rich regions, the locations of the RO data influences these values.

332 So far, this study has focused on nowcast analysis results using only data available  
 333 at the time of the assimilation. It is possible to use methods developed for forecasting  
 334 to further reduce the analysis RMSE by propagating the information from the previous  
 335 assimilation window to the next one by using Gauss-Markov model defined in Equation  
 336 4, because the state starts closer to a realistic ionospheric state. The yellow line in Fig-  
 337 ure 8 shows the RMSEs for the case when the analysis from the previous assimilation  
 338 window was used to update a background for the current window.

339 **4 Discussion**

340 **4.1 Challenges**

341 In the beginning of this paper, the current challenges of the non-parametrized iono-  
 342 spheric DA models were presented. This section discusses how those challenges were ad-  
 343 dressed by ANCHOR.

344 The first challenge was related to the high computational intensity of the non-parametrized  
 345 models, such as IDA4D. Due to the anchor parametrization methods, ANCHOR's co-  
 346 variance matrix is 1347 times smaller than a typical covariance matrix of IDA4D, which  
 347 makes ANCHOR very rapid. On average it takes about 10 sec to produce the analysis  
 348 for one assimilation time frame on a regular PC, while IDA4D takes several minutes when  
 349 configured to use 44 openMP threads on a system with 200 GB of memory.

350 The second challenge was related to information propagation during the assimi-  
 351 lation of non-local data types, such as sTEC from RO data. ANCHOR directly addresses  
 352 this problem through the extraction of the EDP parameters from post-processed RO data,  
 353 thus converting the non-local RO data type into a local measurement. Some may argue  
 354 that an additional data pre-processing, such as Abel inversion for the RO data, intro-  
 355 duces additional errors for the  $NmF2$  parameter. However, as it was demonstrated in  
 356 Figure 7, this approach enables the propagation of the information to all EDP param-  
 357 eters. Additionally, since most DA systems ingest RO sTEC as a relative measurement  
 358 (by the assimilation of the difference between sTEC from a chosen base ray and the other  
 359 rays), it does not guarantee the precise determination of sTEC, nor of  $NmF2$  in addi-  
 360 tion to leaving the other parameters unchanged from the background. In the future, AN-  
 361 CHOR will introduce an advanced RO data pre-processing, by using the information about  
 362 the horizontal density gradients derived from the analysis to improve the currently used

363 Abel inversion technique (by avoiding the symmetry assumption that is used in a stan-  
 364 dard Abel inversion).

365 The third challenge described the presence of unrealistic profiles in IDA4D anal-  
 366 ysis. Since ANCHOR is a parametrized DA model, this issue is fully addressed. Because  
 367 each parameter varies within a physically realistic limit (set by the given model variance)  
 368 it is impossible to obtain a discontinuity, a bite-out, or a fictitious F3 layer in the assimi-  
 369 lation results. This makes ANCHOR a very reliable algorithm, where the output is al-  
 370 ways compatible with the ray tracing applications.

## 371 4.2 ANCHOR in Comparison to Other Parametrized DA Models

372 The parametrization for the ionospheric DA models is a rational approach. Exam-  
 373 ples of two other models that applied the parametrization are IRTAM (Galkin et al., 2020)  
 374 and The Assimilative Canadian High Arctic Ionospheric Model (A-CHAIM) (Reid et al.,  
 375 2023).

376 IRTAM is an operational ionospheric weather model based on ionosonde data from  
 377 the Global Ionosphere Radio Observatory (GIRO) (Reinisch & Galkin, 2011). IRTAM  
 378 adjusts IRI diurnal coefficients by looking at the diurnal time series of the  $NmF2$ ,  $hmF2$ ,  
 379  $B0$  and  $B1$  (the last two are the IRI parameters that describe the F2 bottomside thick-  
 380 ness) parameters at each ionosonde station. Next, IRTAM adjusts global IRI coefficients  
 381 to connect the ionosonde stations and to produce a smooth global distribution of the pa-  
 382 rameters. Since the diurnal distribution of the parameters at fixed locations (like ionosonde  
 383 stations) is the key to the IRTAM, it is challenging to add RO data into the same IR-  
 384 TAM DA formalism, because the locations of RO measurements also change with time.  
 385 Unlike IRTAM, ANCHOR does not perform DA on time series, and therefore is able to  
 386 combine ionosonde and RO data, treating it in a similar way. Comparing the two DA  
 387 systems, IRTAM provides smoothed in time and space global ionospheric background,  
 388 whereas ANCHOR focuses on the nowcasting for each time frame and reveals higher spa-  
 389 cial resolution features in the distribution of the parameters.

390 A-CHAIM applies particle filter DA approach to the series of spherical cap harmonic  
 391 perturbations on the Empirical Canadian High Arctic Ionospheric Model (E-CHAIM).  
 392 Unlike ANCHOR, A-CHAIM parametrizes the ionosphere not only vertically, but also  
 393 horizontally, by finding a set of spherical harmonic coefficients that describe the hori-

394 zontal distribution of the parameters at each time frame. The particle filter DA approach  
 395 makes A-CHAIM somewhat similar to Ensemble KF approach, where a particular re-  
 396 alization of the model is chosen as a combination of all other possible model realizations.  
 397 A-CHAIM is a regional DA model, unlike ANCHOR that produces the analysis on the  
 398 global scale.

## 399 5 Conclusion and Future Directions

400 A new nowcast data assimilation algorithm, ANCHOR, was developed at the U.S.  
 401 Naval Research Laboratory. It extracts ionospheric parameters from RO and ionosonde  
 402 data and assimilates them as point measurements into the maps of the background pa-  
 403 rameters using Kalman Filter approach. ANCHOR is written in Python. It has a high  
 404 computational efficiency, and is guaranteed to provide a reliable analysis suitable for the  
 405 ray tracing applications.

406 There are many avenues for the future development of this DA model. Ground based  
 407 GPS, UV radiance, and HF radar data will be added to the assimilation process. In the  
 408 future ANCHOR will use SAMI3 as a background and will have a sophisticated forecast  
 409 model utilizing the coupling between physics-based SAMI3 model and the ANCHOR now-  
 410 cast. The possibility to assimilate sporadic E data into ANCHOR will be investigated.  
 411 ANCHOR will undergo thorough optimization, stress testing, and validation in future  
 412 efforts, with the goal of providing robust, operational DA support for space weather mod-  
 413 els.

## 414 6 Open Research

- 415 • PyIRI software used as ANCHOR's background model is available to the commu-  
 416 nity at GitHub (Forsythe & Burrell, 2023).
- 417 • Raw COSMIC-2 data is available at <https://data.cosmic.ucar.edu/gnss-ro/cosmic2/provisional/spaceWeather/>.
- 418 • SAMI3 density outputs used as truth here are available at <https://doi.org/10.5281/zenodo.10196102>.
- 419 • Simulated RO and ionosonde parameters are available at <https://doi.org/10.5281/zenodo.10196588>.

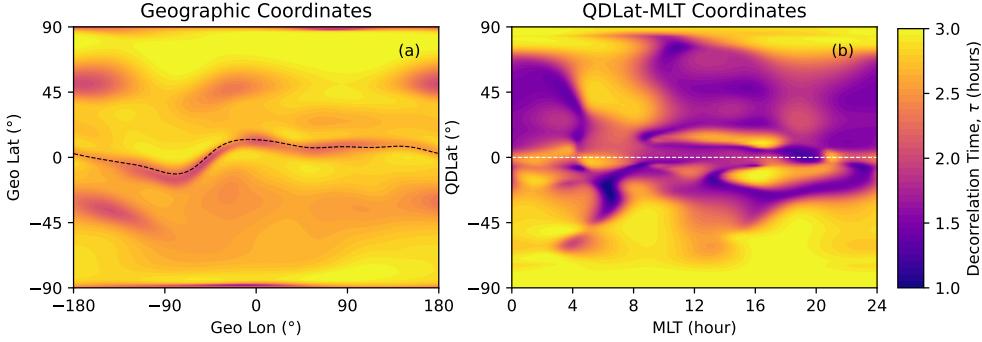
- 423     • ANCHOR assimilation outputs are available at <https://doi.org/10.5281/zenodo.10196650>.
- 424

425     **Acknowledgments**

426     This work is sponsored by the Office of Naval Research.

427     **Appendix A Decorrelation Time**

428     Even though, ANCHOR is developed to be a nowcast assimilation scheme, the ques-  
429     tion about the decorrelation time was investigated. The decorrelation time should in-  
430     dicate the appropriate aging of the data in relation to the forecast background. A long  
431     decorrelation time would indicate slower ionospheric processes and therefore long-lasting  
432     data influence, whereas a short decorrelation time should point to the dynamic ionospheric  
433     region where the data becomes old faster. Figure A1 compares decorrelation time be-  
434     tween geographic and QDLat-MLT coordinate systems. The decorrelation time,  $\tau$  was  
435     calculated from the PyIRI model following the recipe from Forsythe, I. Azeem, et al. (2020).  
436     It is rather counterintuitive, but in the QDLat-MLT system  $\tau$  is shorter in comparison  
437     with the geographic coordinate system, even though the variance is much smaller than  
438     in the geographical system, as was demonstrated in Figure 2. This indicates the pres-  
439     ence of small-amplitude pulsations of the equatorial crests around its MLT and QDLat  
440     position. However, it is questionable whether these pulsations have any physical origin  
441     or whether they arise as an artificial effect that is coming from the URSI coefficients (a  
442     model that determines the global and diurnal distribution of  $NmF2$  in PyIRI) that were  
443     derived for the geographic coordinate system. Further investigation, or even further im-  
444     provement of the climatological coefficients, is required to address this issue.



**Figure A1.** Comparison of the decorrelation time in geographic and QDLat-MLT systems. Geomagnetic equators are shown with dashed lines.

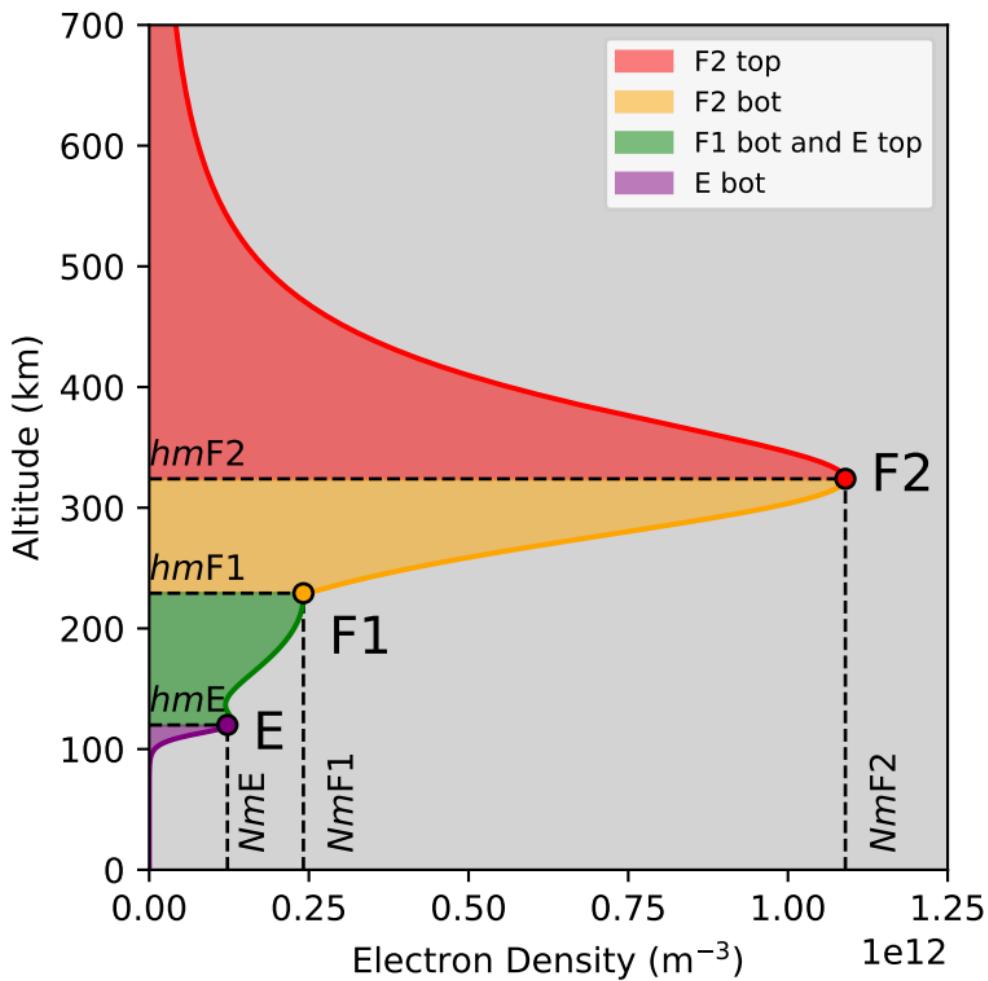
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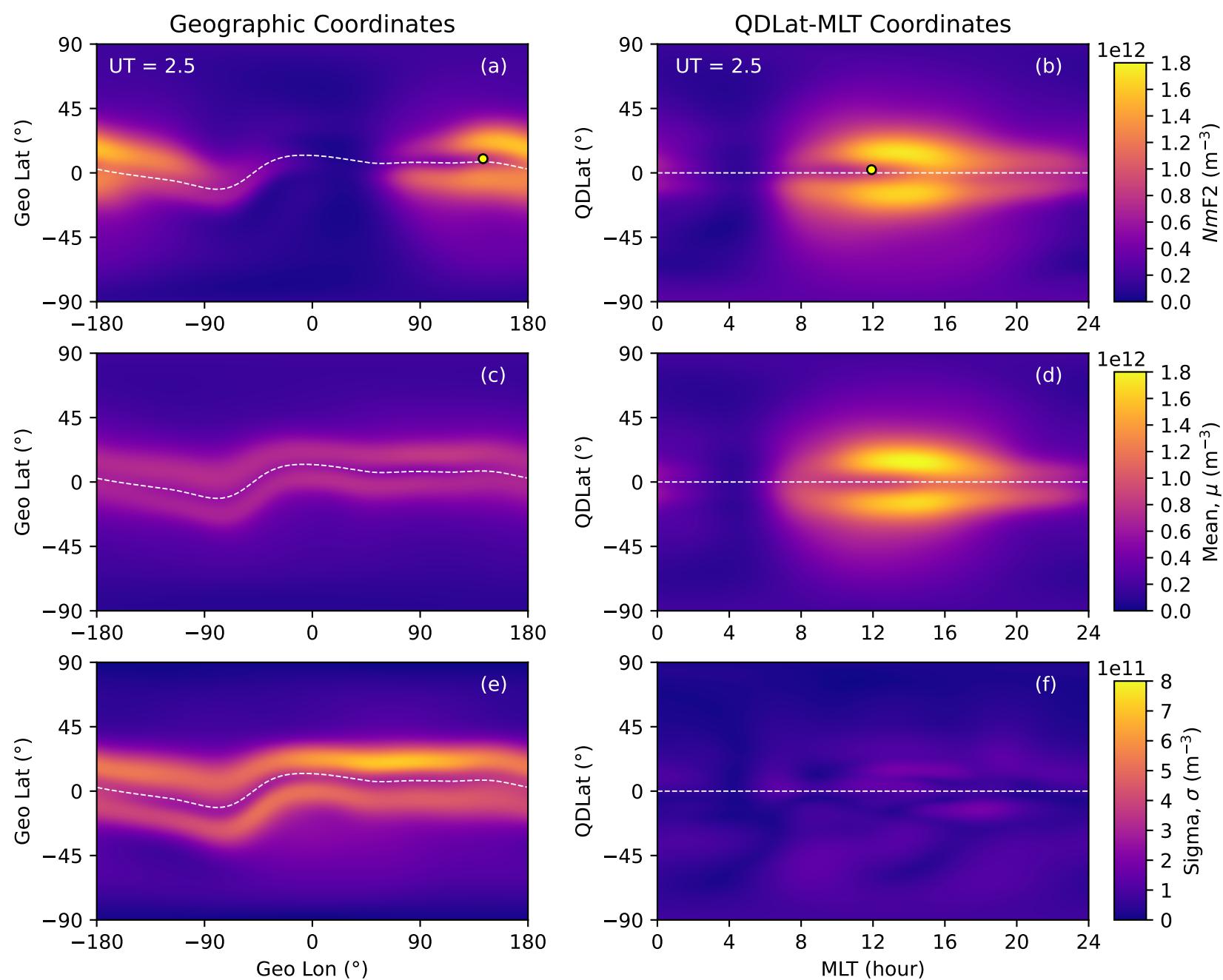
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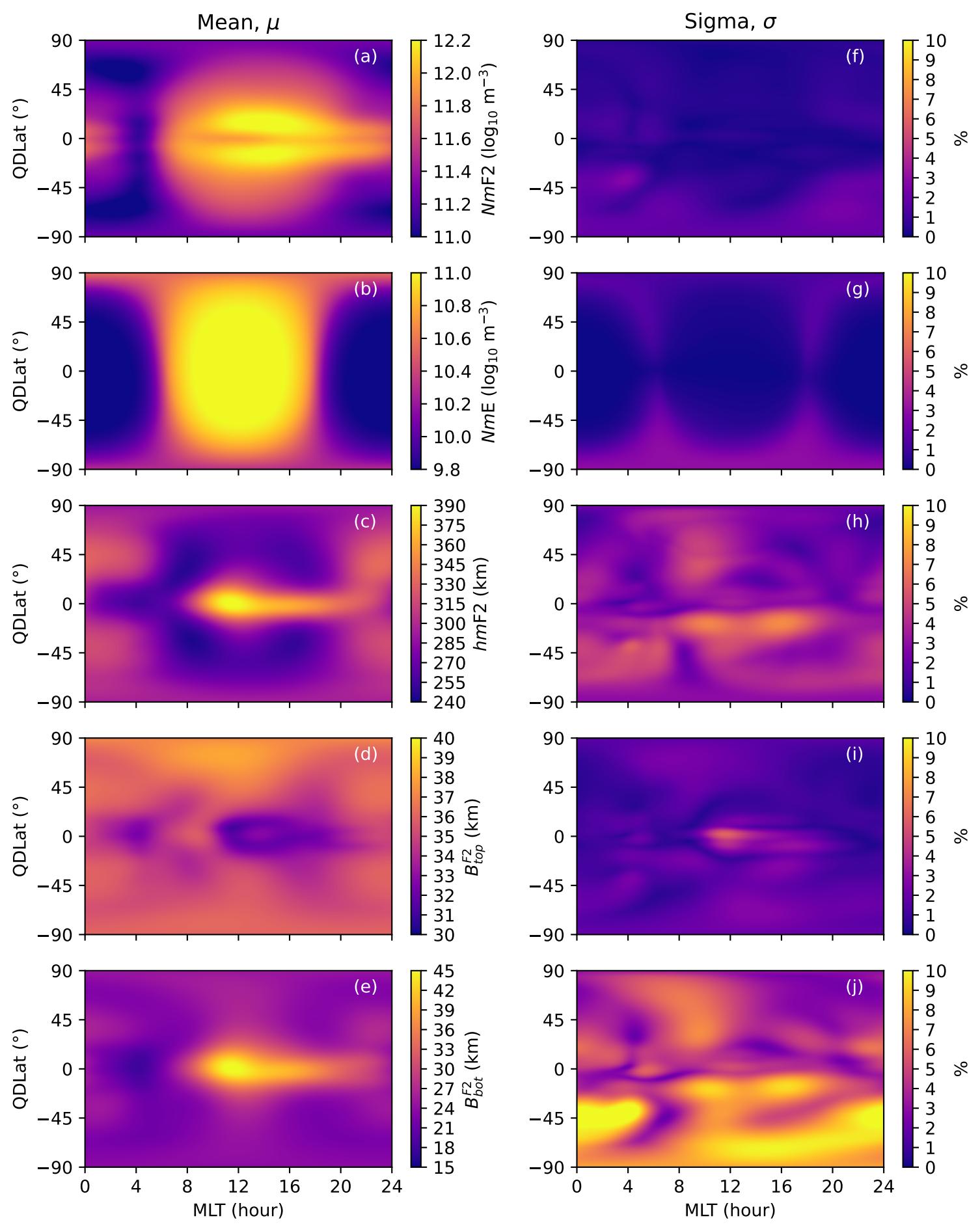
**Figure 1.**



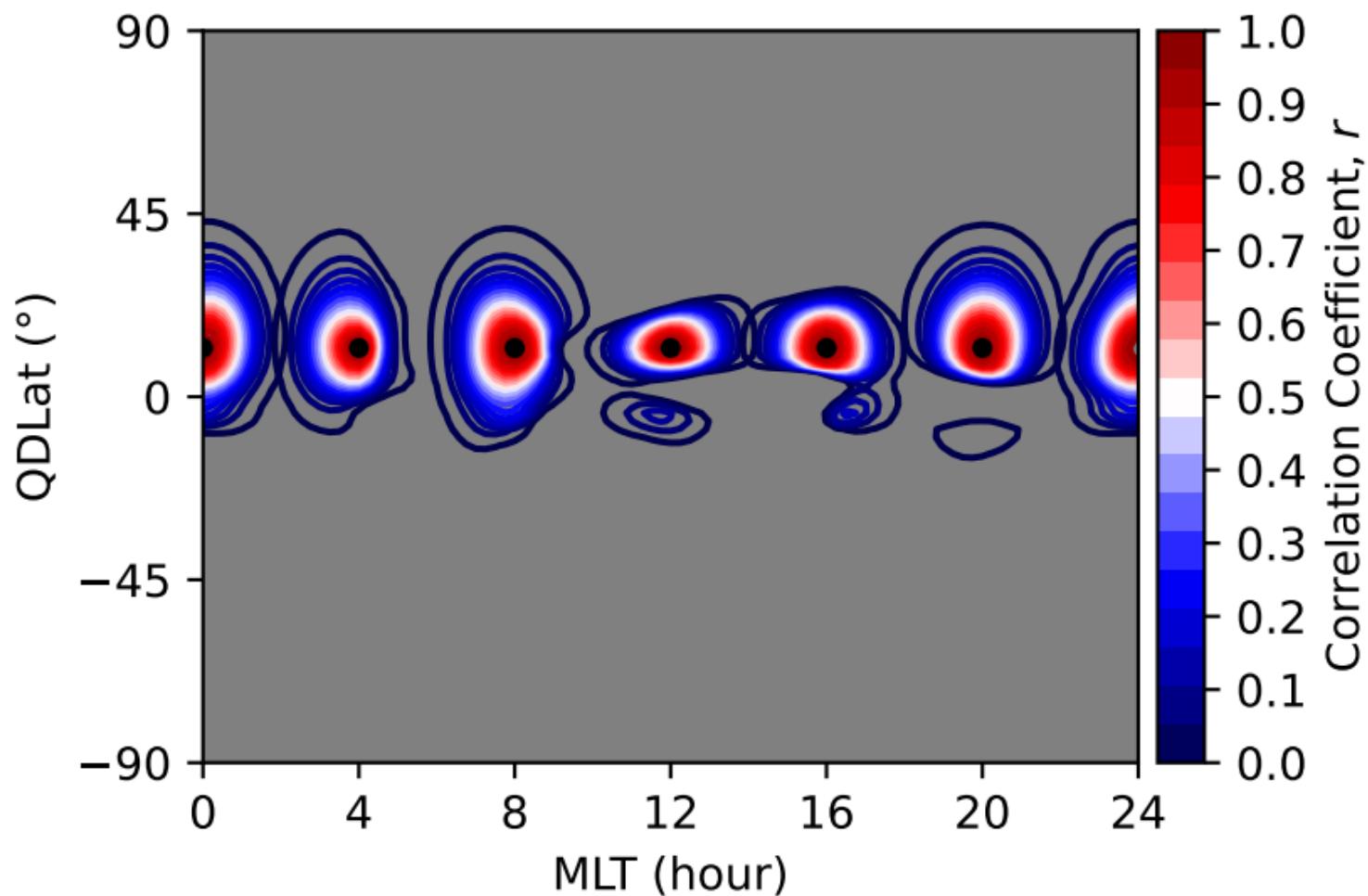
**Figure 2.**



**Figure 3.**

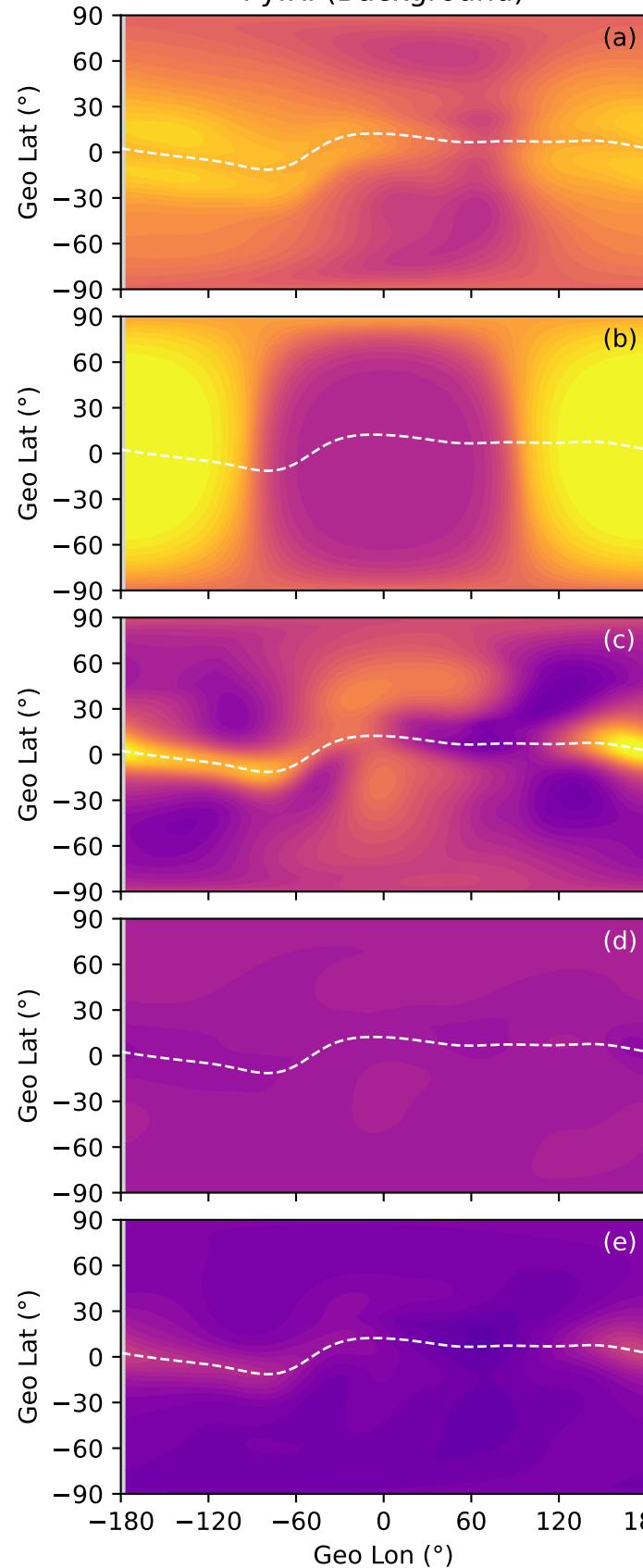


**Figure 4.**

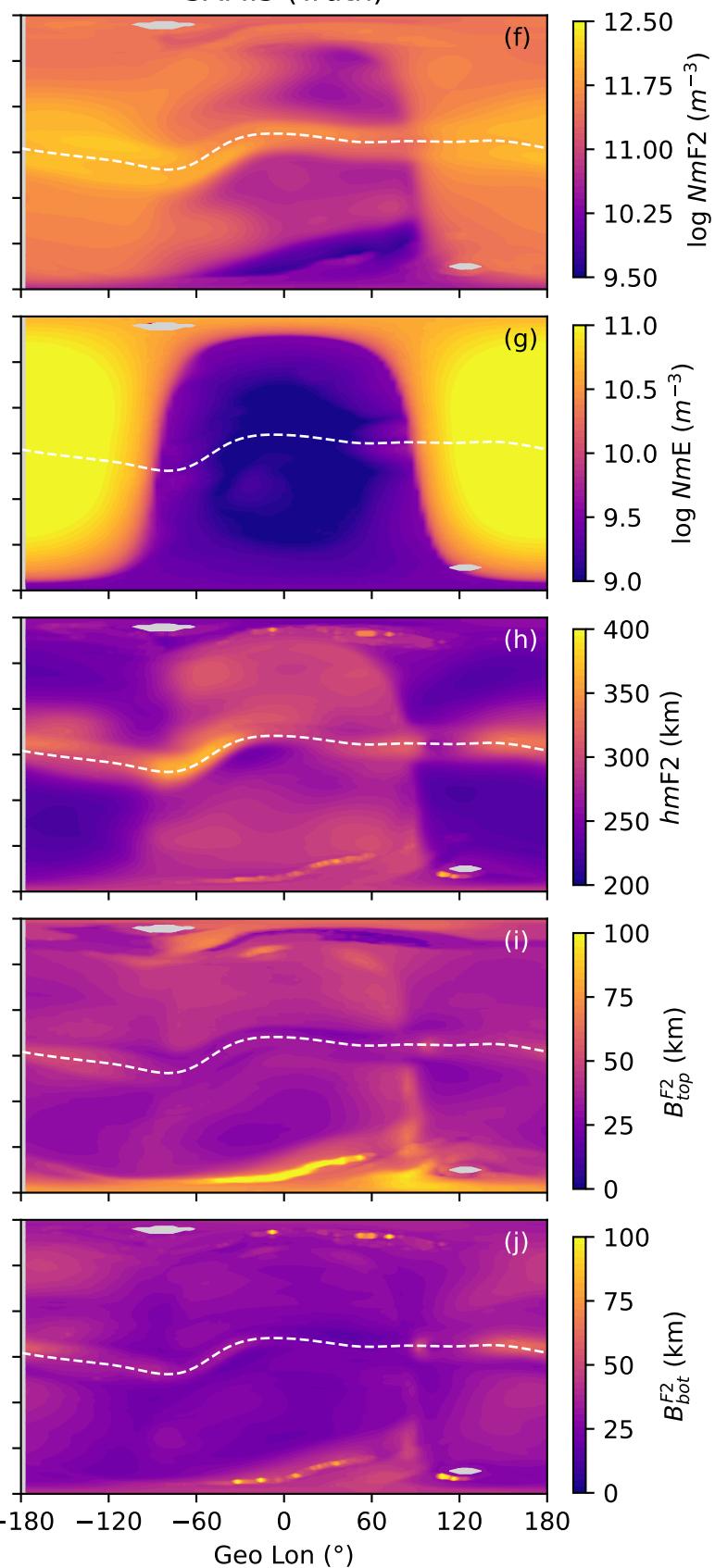


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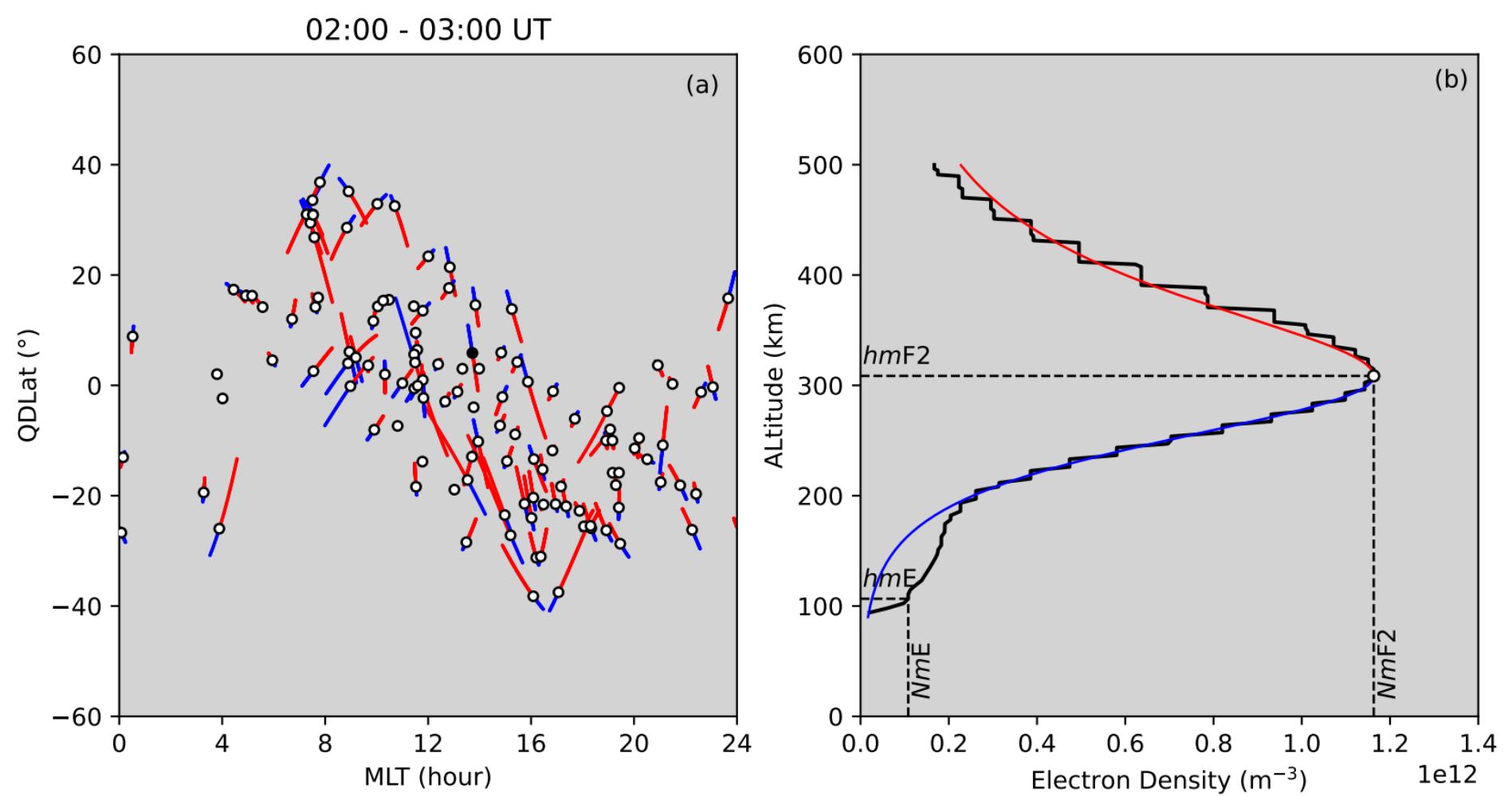
PyIRI (Background)



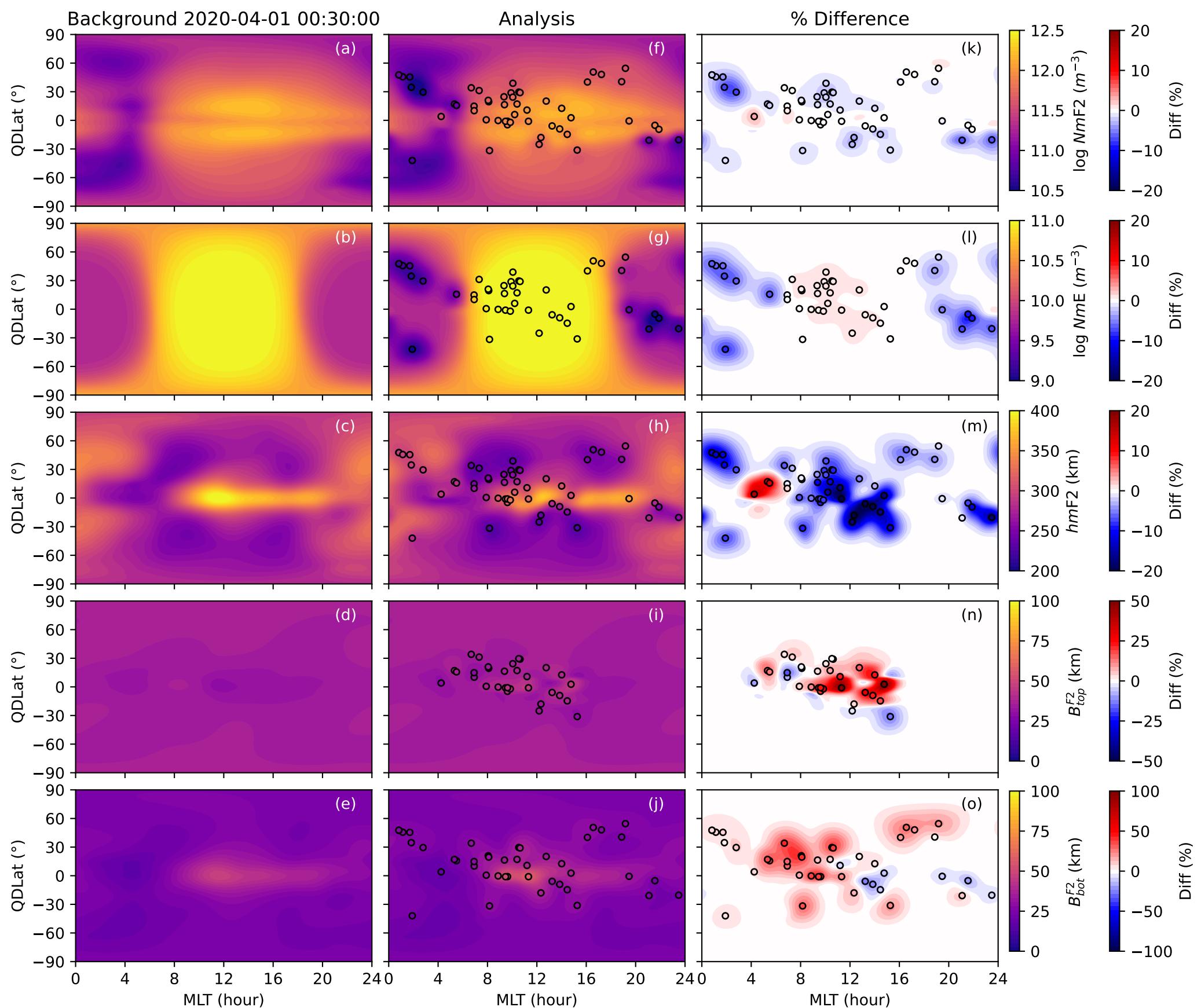
SAMI3 (Truth)



**Figure 6.**

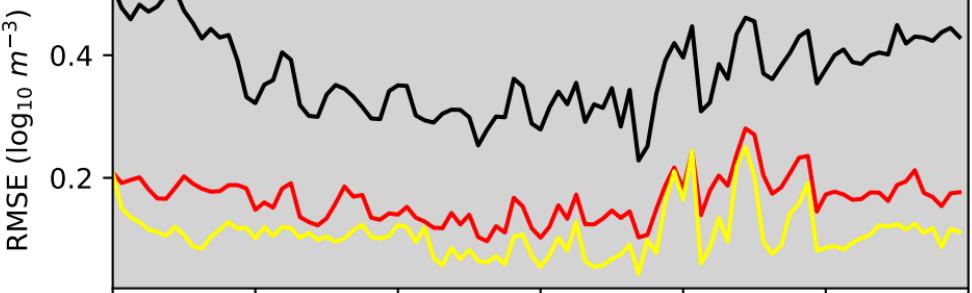


**Figure 7.**

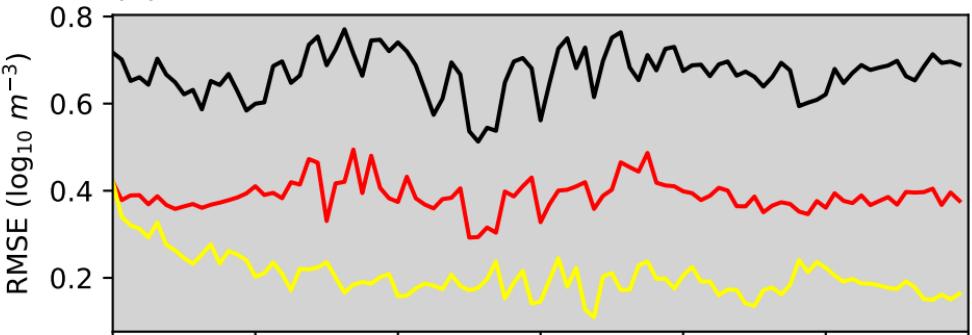


**Figure 8.**

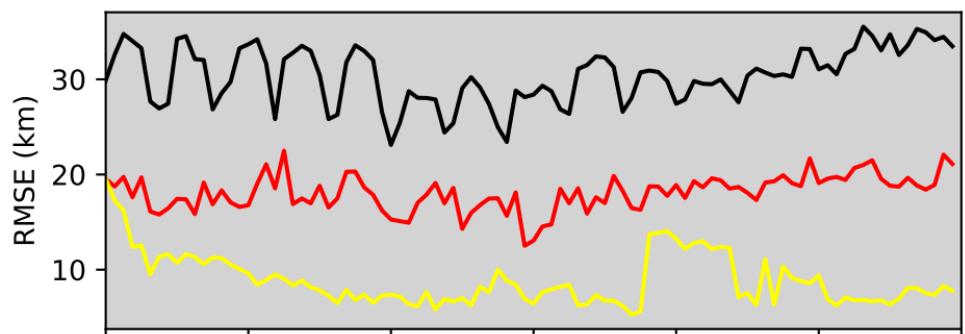
(a)  $NmF2$  — Background — Analysis — Analysis+



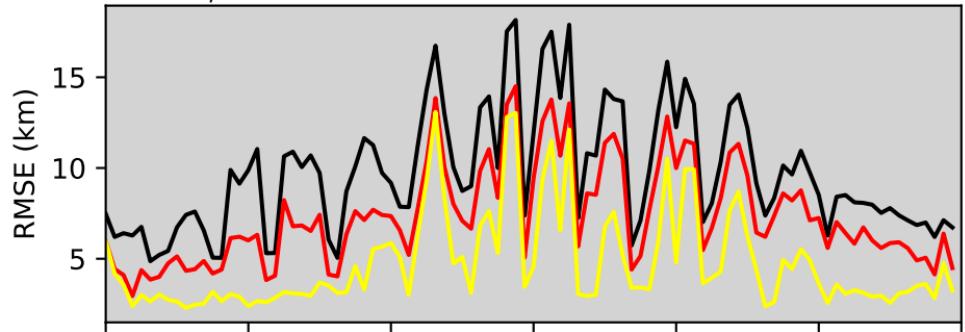
(b)  $NmE$



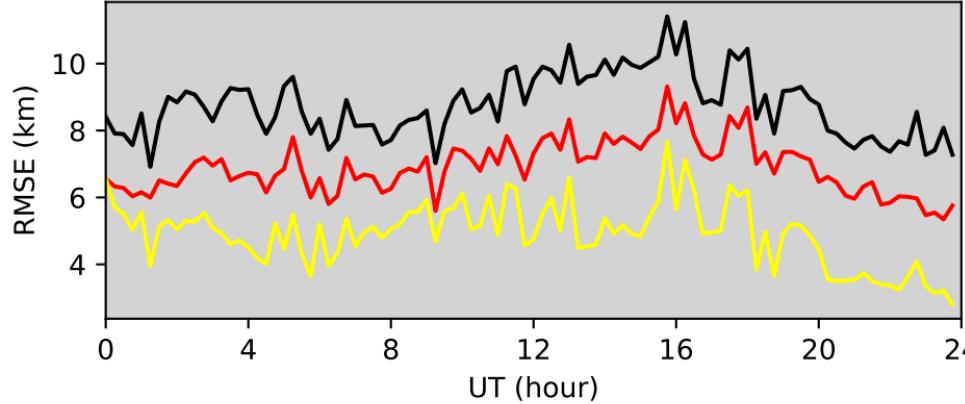
(c)  $hmF2$



(d)  $B_{top}^{F2}$

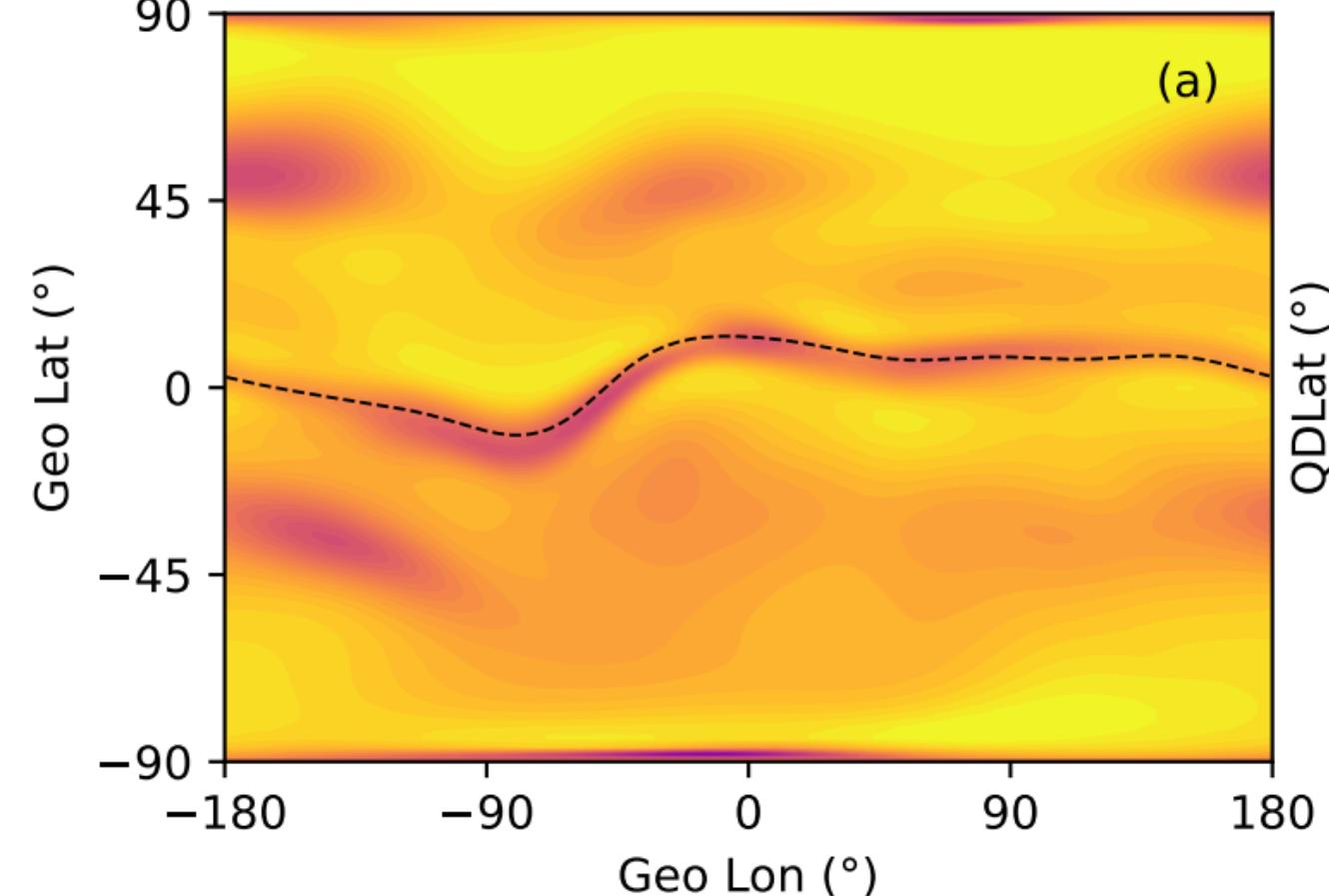


(e)  $B_{bot}^{F2}$

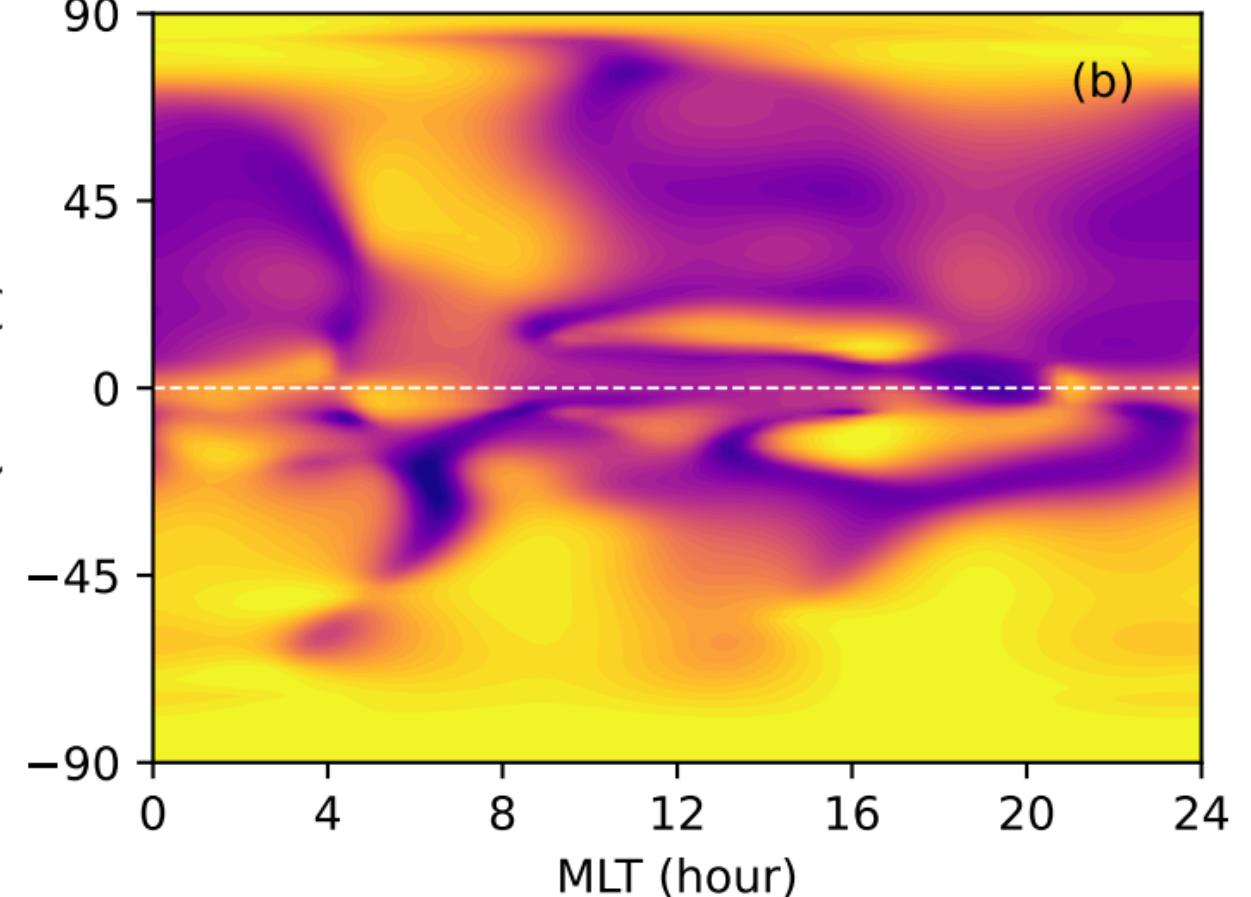


**Figure A1.**

Geographic Coordinates



QDLat-MLT Coordinates



Decorrelation Time,  $\tau$  (hours)

A vertical color bar indicating the value of the decorrelation time  $\tau$  in hours. The scale ranges from 1.0 (dark purple) to 3.0 (bright yellow), with intermediate ticks at 1.5, 2.0, and 2.5.